# EFFECTS OF CHEMICAL REACTION ON DOUBLE DISPERSION HEAT AND MASS TRANSFER IN NON-DARCY POROUS MEDIA

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In this paper, the effects of chemical reaction and double-dispersion on non-Darcy free convection heat and mass transfer from semi-infinite vertical surface in a fluid saturated porous medium are investigated. The Forchheimer extension (nom-Darcy term) is considered in the flow equations, while the chemical reaction power-law term is considered in the concentration equation. The first order chemical reaction (n=1) was used through this study. The thermal and mass diffusivities;  $\alpha_y$  and  $D_y$  are variables defined as  $\alpha_y = \alpha + \gamma d|u|$  and  $D_y = D + \zeta d|u|$ , respectively. The flow, temperature and concentration fields in Darcy and non-Darcy porous media are observed to be governed by complex interactions among the diffusion rate; Le, buoyancy ratio; N, and the parameters  $\gamma$ ,  $\zeta$ ,  $F_0$ ,  $Ra_d$ ,  $\chi$ . Numerical results for the details of the velocity, temperature and concentration profiles as well as heat transfer rate (Nusselt number) and mass transfer rate (Sherwood number) which are shown on graphs have been presented.

**KEYWORDS**: chemical reaction, dispersion; porous medium; heat and mass transfer

# **1. INTRODUCTION**

Chemical reaction effects should be considered in many applications of heat and mass transfer especially, convection in chemical rectors of porous construction. Anjalidevi and Kandasamy [1] studied the effects caused by chemical diffusion mechanics and the inclusion of general chemical reaction of order n on combined forced and natural convection flows over a semi-infinite vertical plate immersed in an ambient fluid. They stated that the presence of pure air or water is impossible in nature and that some foreign mass may be present either naturally or mixed with air or water. Mulolani and Rahman [2] studied laminar natural convection flow over a semi-infinite vertical plate under the assumption that the concentration of species along the plate

NOMENCLATURE				
$\mathbf{F}_0$	Structural and thermo	β	Thermal expansion coefficient	
	physical parameter	$\beta^*$	Solutal expansion coefficient	
G	Gravitational accleration	χ	Non-dimensional chemical reaction-porous medium parameter	
Gc	Modified Grashof number,	φ	Dimensionless Concentration	
	$\left[=\beta^*g(C_w - C_{\infty}) x^3/v^2\right]$	γ	Mechanical thermal dispersion coefficient	
$\mathbf{J}_{\mathbf{w}}$	Local mass flux	η	Similarity space variable	
К	Molecular thermal conductivity	λ	Non-dimensional chemical reaction parameter	
K	Permeability of the porous medium	μ	Fluid dynamic viscosity	
$\mathbf{K}_0$	Chemical reaction parameter	ν	Fluid kinematic viscosity	
k <sub>d</sub>	Dispersion thermal conductivity	θ	Dimensionless temperature	
ke	Effective thermal conductivity	ρ	Fluid density	
Le	Lewis number, $[= \alpha/D]$	Ψ	Stream function	
N	Order of reaction	ζ	Mechanical solutal dispersion coefficient	
Ν	Buoyancy ratio			
Nu <sub>x</sub>	Local Nusselt number	Sub	Subscripts	
Р	Pressure			
Pr	Prandtl number, $[= \nu/\alpha]$	DI	Pore diameter	
Q	Heat transfer rate	X, I y	n the direction of x and y	
Ra	Rayleigh number	W S	Surface conditions	
Re <sub>x</sub>	Local Reynolds number	∞ (	Conditions away from the surface	
Sc	Schmidt number, $[= \nu/D]$			
$\mathbf{Sh}_{\mathbf{x}}$	Local Sherwood number	Sup	erscripts	
Т	Temperature	′ I	Derivative with respect to $\eta$	

follows some algebraic law with respect to chemical reaction. They obtained similarity solutions for different orders of reaction; n and Schmidt number; Sc. At constant temperature, the rate of chemical reaction depends on the concentrations of reactants and products [3], it may also depends on catalysts and inhibitors. The rates of chemical reactions are obtained from measurements of concentration as a function of time. The rate of reaction depends primarily on the order of the reaction with respect to reactant A. Prasad et al. [4] studied the influence reaction rate on the transfer of chemically reactive species in the laminar visco-elastic fluid flow immersed in porous medium over a stretching sheet. They concluded that the effect of chemical reaction is to reduce the thickness of concentration boundary layer and to increase the mass transfer rate

from the sheet to the surrounding fluid and that this effect is more effective for zero and first order reactions than second and third order reactions.

Study of the thermal dispersion effects become prevalent in the porous media flow region. The thermal and solutal dispersion effects become more important when the inertial effects are prevalent. Fried and Combarnous [5] proposed a linear function to express the thermal dispersion. Also, a linear dispersion model takes the porosity of the porous medium into account is used for free convection in a horizontal layer heated from below was introduced by Georgiadis and Catton [6]. Cheng [7] and Plumb [8] gave another model for flow and heat transfer in porous media by taking thermal dispersion effects into consideration. An analysis of thermal dispersion effect on vertical plate natural convection in porous media is presented by Hong and Tien [9]. Lai and Kulacki [10] investigated thermal dispersion effects of thermal dispersion and lateral mass flux on non-Darcy natural convection over a vertical flat plate in a fluid saturated porous medium were studied by Murthy and Singh [11].

The complexity of the flow increases when higher order effects like thermal and solutal dispersion are considered in the medium. Karimi-Fard et al. [12] presented a numerical study of double-diffusive free convection heat and mass transfer in a square cavity filled with a porous medium. Began [13] has been analyzed the effect of solutal and thermal dispersion in Darcian porous medium. The double dispersion phenomenon in a free convection boundary layer adjacent to a vertical wall in a Darcian porous medium, using scale analysis arguments, was investigated by Telles and Trevisan [14]. Effects of double dispersion on mixed convection heat and mass transfer in non-Darcy porous medium has been investigated by Murthy [15]. Recently, El-Amin [16] has studied the problem of double-dispersion on non-Darcy free convection heat and mass transfer over a vertical flat plate in a fluid saturated porous medium.

The present investigation is devoted to study the combined effect of chemical reaction and solutal/thermal dispersion on Forchheimer natural convection heat and mass transfer over a vertical flat plate in a fluid saturated porous medium. The Forchheimer extension (non-Darcy term) is considered in the flow equations, while the chemical reaction power-law term is considered in the concentration equation. The first order chemical reaction (n=1) was used through this study. The thermal and mass diffusivities are defined in terms of the molecular thermal and solutal diffusivities respectively. The flow, temperature and concentration fields in Darcy and non-Darcy porous media are observed to be governed by complex interactions among the diffusion rate; Le, buoyancy ratio; N, and the parameters  $\gamma$ ,  $\zeta$ ,  $F_0$ ,  $Ra_d$ ,  $\chi$ . The wall temperature and the wall concentration distribution are assumed to be uniform.

#### 2. FORMULATION OF THE PROBLEM

In the present study, the effect of chemical reaction; of order n, on non-Darcy natural convection heat and mass transfer over a semi -infinite vertical surface in a fluid saturated porous medium is considered. The x-axis is taken along the plate and the y-axis is normal to it. The wall is maintained at constant temperature and concentration,  $T_w$  and  $C_w$ , respectively, and these values are assumed to be greater than the ambient

temperature and concentration,  $T_{\infty}$  and  $C_{\infty}$ , respectively. The governing equations for this problem; [2] and [16] are given by:

#### Mass:

$$\frac{\partial \mathbf{u}}{\partial \mathbf{x}} + \frac{\partial \mathbf{v}}{\partial \mathbf{y}} = 0 \tag{1}$$

Momentum:

$$u + \frac{c\sqrt{K}}{v}uq = -\frac{K}{\mu}(\frac{\partial p}{\partial x} + \rho g)$$
(2)

$$v + \frac{c\sqrt{K}}{v}vq = -\frac{K}{\mu}(\frac{\partial p}{\partial y})$$
(3)

**Energy:** 

$$u\frac{\partial T}{\partial x} + v\frac{\partial T}{\partial y} = \frac{\partial}{\partial x}(\alpha_x\frac{\partial T}{\partial x}) + \frac{\partial}{\partial y}(\alpha_y\frac{\partial T}{\partial y})$$
(4)

#### **Concentration:**

$$u\frac{\partial C}{\partial x} + v\frac{\partial C}{\partial y} = \frac{\partial}{\partial x}(D_x\frac{\partial C}{\partial x}) + \frac{\partial}{\partial y}(D_y\frac{\partial C}{\partial y}) - K_0(C - C_{\infty})^n$$
(5)

**Density:** 

$$\rho = \rho_{\infty} [1 - \beta (T - T_{\infty}) - \beta^* (C - C_{\infty})]$$
<sup>(6)</sup>

along with the boundary conditions

$$y = 0: v = 0, \quad T_w = \text{const.}, \quad C_w = \text{const.}, y \to \infty: u = 0, \quad T \to T_{\infty}, \quad C \to C_{\infty}.$$
(7)

The chemical reaction effect is acted by the last term of Eq. (5), where, the power n is the order of reaction and  $K_0$  is the chemical reaction parameter. The normal component of the velocity near the boundary is small compared with the other component of the velocity and the derivatives of any quantity in the normal direction are large compared with derivatives of the quantity in direction of the wall. Under these assumptions, Eq. (1) remains the same, while Eqs. (2)-(6) become

$$u + \frac{c\sqrt{K}}{v}u|u| = -\frac{K}{\mu}(\frac{\partial p}{\partial x} + \rho g)$$
(8)

$$\frac{\partial \mathbf{p}}{\partial \mathbf{y}} = 0 \tag{9}$$

$$u\frac{\partial T}{\partial x} + v\frac{\partial T}{\partial y} = \frac{\partial}{\partial y}(\alpha_y \frac{\partial T}{\partial y})$$
(10)

$$u\frac{\partial C}{\partial x} + v\frac{\partial C}{\partial y} = \frac{\partial}{\partial y}(D_y\frac{\partial C}{\partial y}) - K_0(C - C_{\infty})^n$$
(11)

Following Telles and Trevisan [14], the quantities of  $\alpha_y$  and  $D_y$  are variables and are defined as  $\alpha_y = \alpha + \gamma d|u|$  and  $D_y = D + \zeta d|u|$ , where,  $\alpha$  and D are the molecular thermal and solutal diffusivities, respectively, whereas  $\gamma d|u|$  and  $\zeta d|u|$  represent dispersion thermal and solutal diffusivities, respectively. This model for thermal dispersion has been used extensively by researchers like Cheng [7], Plumb [8], Hong and Tien [9], Lai and Kulacki [10] and Murthy and Singh [11] in studies of convective heat transfer in non-Darcy porous media.

Having invoked the Boussinesq approximations, with substituting Eq. (6) into Eqs. (8) and (9), eliminating the pressure and the velocity components u and v can be written in terms of stream function  $\psi$  as:  $u = \partial \psi / \partial y$  and  $v = -\partial \psi / \partial x$ , we obtain:

$$\frac{\partial^2 \psi}{\partial y^2} + \frac{c\sqrt{K}}{\nu} \frac{\partial}{\partial y} (\frac{\partial \psi}{\partial y})^2 = \left(\frac{Kg\beta}{\mu} \frac{\partial T}{\partial y} + \frac{Kg\beta^*}{\mu} \frac{\partial C}{\partial y}\right) \rho$$
(12)

$$\frac{\partial \psi}{\partial y} \frac{\partial T}{\partial x} - \frac{\partial \psi}{\partial x} \frac{\partial T}{\partial y} = \frac{\partial}{\partial y} \left[ (\alpha + \gamma d \frac{\partial \psi}{\partial y}) \frac{\partial T}{\partial y} \right]$$
(13)

$$\frac{\partial \psi}{\partial y} \frac{\partial C}{\partial x} - \frac{\partial \psi}{\partial x} \frac{\partial C}{\partial y} = \frac{\partial}{\partial y} \left[ (D + \zeta d \frac{\partial \psi}{\partial y}) \frac{\partial C}{\partial y} \right] - K_0 (C_w - C_w)^n \phi^n$$
(14)

Introducing the similarity variable and similarity profiles; [16]:

$$\eta = \operatorname{Ra}_{x}^{1/2} \frac{y}{x}, f(\eta) = \frac{\psi}{\alpha \operatorname{Ra}_{x}^{1/2}}, \theta(\eta) = \frac{T - T_{\infty}}{T_{w} - T_{\infty}}, \phi(\eta) = \frac{C - C_{\infty}}{C_{w} - C_{\infty}}$$
(15)

where  $Ra_x$  is the modified local Rayleigh number; [16],  $Ra_x = Kg\beta(T_w - T_{\infty})x/\alpha v$ . The problem statement then becomes:

$$\mathbf{f}'' + 2\mathbf{F}_0 \mathbf{Ra}_{\mathbf{d}} \mathbf{f}' \mathbf{f}'' = \mathbf{\theta}' + \mathbf{N} \mathbf{\phi}' \tag{16}$$

$$\theta'' + \frac{1}{2}f\theta' + \gamma Ra_{d}(f'\theta'' + f''\theta') = 0$$
<sup>(17)</sup>

$$\phi'' + \frac{1}{2} \operatorname{Lef} \phi' + \zeta \operatorname{Le.Ra}_{d} (f' \phi'' + f'' \phi') - \operatorname{Sc} \lambda \frac{\operatorname{Gc}}{\operatorname{Re}_{x}^{2}} \phi^{n} = 0$$
(18)

With analogy to [2] and [18] Gc is the modified Grashof number,  $Re_x$  is local Reynolds number, Sc and  $\lambda$  are Schmidt number and non-dimensional chemical reaction parameter defined as

$$Gc = \beta^{\bullet}g(C_{w} - C_{\infty})^{2} x^{3} / v^{2}, Re_{x} = u_{r}x/v$$

$$\lambda = \frac{K_{0}\alpha d}{Kg\beta^{*}}(C_{w} - C_{\infty})^{n-3}$$
(19)

and

where the diffusivity ratio (Lewis number), Le is the ratio of Schmidt number and Prandtl number, and  $u_r$  is reference velocity; [17] defined as,

$$u_{r} = \sqrt{g\beta d(T_{w} - T_{\infty})}$$
<sup>(20)</sup>

Equation (18) can be rewritten in the following form

$$\phi'' + \frac{1}{2} \operatorname{Lef} \phi' + \zeta \operatorname{Le.Ra}_{d} (f' \phi'' + f'' \phi') - \chi \phi^{n} = 0$$
(18)

With analogy to [4] and [18], the non-dimensional chemical reaction -porous medium parameter;  $\chi$  is defined as

$$\chi = \mathrm{Sc}\lambda \frac{\mathrm{Gc}}{\mathrm{Re}_{\mathrm{x}}^2} \tag{21}$$

The boundary conditions become

$$f(0) = 0, \theta(0) = \phi(0) = 1, \ f'(\infty) = \theta(\infty) = \phi(\infty) = 0$$
(22)

As mentioned in [16], the parameter  $F_0 = c\sqrt{K\alpha}/\nu d$  represents the structural and thermo physical properties of the porous medium,  $Ra_d = Kg\beta(T_w - T_w)d/\alpha\nu$  is the pore diameter dependent Rayleigh number which describes the relative intensity of the buoyancy force, such that *d* is the pore diameter, and  $N = \beta * (C_w - C_w)/\beta(T_w - T_w)$  is the buoyancy ratio.

It is noteworthy that  $F_0 = 0$  corresponds to the Darcian free convection,  $\gamma = 0$  represents the case where the thermal dispersion effect is neglected and  $\zeta = 0$  represents the case where the solutal dispersion effect is neglected. In Equation (16), N > 0 indicates the aiding buoyancy and N < 0 indicates the opposing buoyancy.

From the definition of the stream function, the velocity components; [16] become

$$u = \frac{\alpha}{x} Ra_{x} f', \ v = -\frac{\alpha}{2x} Ra_{x}^{1/2} [f - \eta f']$$
(23)

The local heat transfer rate which is the primary interest of the study is given by; [16]

$$q_{w} = -k_{e} \frac{\partial T}{\partial y}\Big|_{y=0} = -(k+k_{d}) \frac{\partial T}{\partial y}\Big|_{y=0}$$
(24)

where  $k_e$  is the effective thermal conductivity of the porous medium which is the sum of the molecular thermal conductivity k and the dispersion thermal conductivity  $k_d$ .

Together with the definition of the local Nusselt number Nu<sub>x</sub>; [16]

$$Nu_{x} = \frac{q_{w}}{T_{w} - T_{\infty}} \frac{x}{k_{e}}$$
(25)

Hence,

$$Nu_{x}Ra_{x}^{-1/2} = -[1 + \gamma Ra_{d}F'(0)]\theta'(0)$$
(26)

Also, the local mass flux is given by

$$j_{w} = -D_{y} \frac{\partial C}{\partial y} \bigg|_{y=0}$$
(27)

Therefore, the Sherwood number; [16] is defined by

$$Sh_{x}Ra_{x}^{-1/2} = -[1 + \zeta Ra_{d}F'(0)]\phi'(0)$$
(28)

#### 3. RESULTS AND DISCUSSION

In order to get the physical insight, the system of ordinary differential equations (16)-(18) along with the boundary conditions (22), are integrated numerically by means of the fourth-order Runge-Kutta method with shooting technique. The step size  $\Delta \eta = 0.05$  is used while obtaining the numerical solution with  $\eta_{max} = 4$  and five-decimal accuracy as the criterion for convergence. Numerical computations are carried out for  $0 \le F_0 \le 0.5$ ,  $0 \le Le \le 4.0$ ,  $0.5 \le Ra_d \le 1.0$ ,  $0 \le \gamma, \zeta \le 0.5$ , and  $0.00 \le \chi \le 0.08$  for the two cases aiding buoyancy (N>0) and opposing buoyancy (N<0) and first order of chemical reaction (n = 1).

The behaviors of velocity and temperature profiles for different values of the non-Darcy parameter;  $F_0$  are shown in **Figs. 1** and **2**, respectively. It is clear from **Fig. 1** that the non-Darcy parameter reduces the velocity close to the wall while it has no effect far from it. **Figure 2** illustrates that there is slight increase in the temperature as a result of increasing the non-Darcy parameter can be shown in **Fig. 2**.



**Figure 1:** Variation of dimensionless velocity ; f' with similarity space variable;  $\eta$  for different F<sub>0</sub> ( Le  $\Box$ = 0.5, Ra<sub>d</sub> = 0.7, N = -0.1,  $\gamma$  =  $\zeta$  = 0.0,  $\chi$  = 0.02)

**Figure 2:** Variation of dimensionless temperature;  $\theta$  with similarity space variable;  $\eta$  for different F<sub>0</sub> (Le  $\Box$  = 0.5, Ra<sub>d</sub> = 0.7, N = -0.1,  $\gamma = \zeta = 0.0$ ,  $\chi = 0.02$ ).

Effects of Lewis number on Nusselt number and Sherwood number for various  $F_0$  are shown in **Figs. 3** and **4**, respectively. It is interesting to note, from **Fig. 3**, that the parameter  $F_0$  has a clear effect on the heat transfer rate while Le has a small effect on it. The opposite is true for the mass transfer rate as illustrated in **Fig. 4**. Moreover, increasing of the non-Darcy parameter;  $F_0$  causes a reduction in the heat transfer rate, while, Lewis number Le enhances it slightly. On the other side, there is an evident effect (reducing) of the parameter  $F_0$  on the mass transfer rate (Sherwood number) with increasing of the Le values.



**Figure 3:** Effect of Lewis number on Nusselt number for various  $F_0$  ( $\chi = 0.02$ , N = -0.1, Ra<sub>d</sub> = 0.7,  $\gamma = \zeta = 0.0$ )



**Figure 4:** Effect of Lewis number on Sherwood number for various  $F_0$  ( $\chi = 0.02$ , N = -0.1, Ra<sub>d</sub> = 0.7,  $\gamma = \zeta = 0.0$ )

**Figures 5** and **6** illustrate the effects of Le on velocity and concentration profiles, respectively. The sensitivity of concentration distribution for Lewis number Le is higher than that of the velocity function. By looking for the non-dimensional concentration equation, we find that Le appears in most of equation terms, which may interprets this behavior mathematically. By the way, as Le increases the velocity profile increases while the concentration distribution decreases.



**Figure 5:** Effect of Le on velocity variation ( $F_0 \square = 0.3$ ,  $Ra_d = 0.7$ , N = -0.1,  $\gamma = \zeta = 0.0$ ,  $\chi = 0.02$ ).

**Figure 6:** Effect of Le on concentration variation ( $F_0 \square = 0.3$ ,  $Ra_d = 0.7$ , N = -0.1,  $\gamma = \zeta = 0.0$ ,  $\chi = 0.02$ ).

**Figures 7-9** show the velocity, temperature and concentration, respectively, as functions of the similarity variable  $\eta$  for various values of the buoyancy parameter; N. It is observed, from these figures, that the buoyancy enhances the velocity while it

reduces both of temperature and concentration profiles. Variations of both Nusselt number and Sherwood number with the buoyancy ratio N are plotted as a function of Le in **Figs. 10** and **11**, respectively. It can be seen that the buoyancy ratio N enhances both of heat and mass transfer rates.

The behavior of the velocity profiles according to the variation of the parameter  $Ra_d$  is shown in **Fig. 12**. It is clear from this that the parameter  $Ra_d$  reduces the velocity close to the wall while it has no effect far from it.

Effects of Lewis number on Nusselt number and Sherwood number for various  $Ra_d$  are shown in **Figs. 13** and **14**, respectively. It is interesting to note that, increasing of the parameter  $Ra_d$  causes a reduction in both of the heat transfer rate and mass transfer rate.





**Figure 7:** Variation of dimensionless velocity ; f ' with similarity space variable;  $\eta$  for different N (L e  $\Box$ = 0.5, F<sub>0</sub> = 0.3, Ra<sub>d</sub> = 0.7,  $\gamma$  =  $\zeta$  = 0.0,  $\chi$  = 0.02).

**Figure 8:** Variation of dimensionless temperature;  $\theta$  with similarity space variable;  $\eta$  for different N (Le = 0.5, F<sub>0</sub> = 0.3, Ra<sub>d</sub> = 0.7,  $\gamma = \zeta = 0.0, \chi = 0.02$ )



N = -0.10.54 N= 0.0 V= 0 1 0.52 N= 0 3  $Nu_x/(Ra_x^{\Lambda}0.5)$ N = 0.50.50 0.48 0.46 0.44 0.42 0.0 1.0 2.0 3.0 4.0 Le

**Figure 9:** Variation of dimensionless concentration;  $\phi$  with similarity space variable;  $\eta$  for different N (Le = 0.5, F<sub>0</sub> =

**Figure 10:** Effect of Lewis number on Nusselt number for various N ( $\chi = 0.02$ ,  $F_0 = 0.3$ ,  $Ra_d = 0.7$ ,  $\gamma = \zeta = 0.0$ ).



0.3,  $Ra_d = 0.7$ ,  $\gamma = \zeta = 0.0$ ,  $\chi = 0.02$ ).

**Figure 11:** Effect of Lewis number on Sherwood number for various N ( $\chi = 0.02$ , F<sub>0</sub> = 0.3, Ra<sub>d</sub> = 0.7,  $\gamma = \zeta = 0.0$ )



**Figure 12:** Velocity variation at different values of Ra<sub>d</sub> (Le  $\Box$  = 0.5,F0 = 0.3, N = -0.1,  $\gamma$  =  $\zeta$  = 0.0,  $\chi$  = 0.02)



Figure 13 Variation of Nusselt number with Lewis number for various  $Ra_d$  ( $\chi = 0.02$ , N = -0.1, F<sub>0</sub> = 0.3,  $\gamma = \zeta = 0.0$ )



Figure 14 Variation of Sherwood number with Lewis number for various  $Ra_d$  ( $\chi = 0.02$ , N = -0.1, F<sub>0</sub> = 0.3,  $\gamma = \zeta = 0.0$ )

Effects of the chemical reaction parameter  $\chi$  on the concentration parameter as a function of the boundary layer thickness are plotted in **Fig. 15**. The figure indicates that as the chemical parameter increases the concentration profiles decreases. As it may be expected, the chemical reaction parameter enhances the mass transfer rate in terms of Sherwood number as shown in **Fig. 16**. It is worth mentioning that the effects of chemical reaction on velocity and temperature profiles as well as heat transfer rate are negligible.

**Figures 17** and **18** illustrate, respectively, the effect of Lewis number on Nusselt number and Sherwood number for various  $\zeta \Box$ . The parameter  $\zeta$  reduces the heat transfer rate especially with high Le values. In the case of mass transfer rate in

terms of Sherwood number, the parameter  $\zeta$  increases the mass transfer rate with small values of Le and the opposite is true for high values of Le. Figure 19 indicates that the thermal dispersion parameter enhances the heat transfer rate.





**Figure 15:** Variation of dimensionless concentration;  $\phi$  with similarity space variable;  $\eta$  for different  $\chi$  (Le = 0.5, F<sub>0</sub> = 0.3, Ra<sub>d</sub> = 0.7,  $\gamma = \zeta = 0.0$ , N = -0.1).

**Figure 16:** Effect of Lewis number on Sherwood number for various  $\chi \square (F_0 = 0.3, N = -0.1, Ra_d = 0.7, \gamma = \zeta = 0.0).$ 



**Figure 17:** Variation of Nusselt number with Lewis number for various  $\zeta \Box$  ( $\chi$ =0.02. Ra<sub>d</sub> = 0.7, F<sub>0</sub> = 0.3, N = - 0.1,  $\Box \gamma$  = 0.0 $\Box$ )

**Figure 18:** Effect of Lewis number on Sherwood number for various  $\zeta \Box$  ( $\chi$  =0.02. Ra<sub>d</sub> = 0.7, F<sub>0</sub> = 0.3, N = -0.1,  $\Box \Box \gamma = 0.0\Box$ )

# 4. CONCLUSIONS

The effects of chemical reaction and double-dispersion on non-Darcy free convection heat and mass transfer from semi-infinite vertical surface in a fluid saturated porous medium are investigated in this article. The flow, temperature and concentration fields are observed to be governed by complex interactions among the diffusion rate Le, buoyancy ratio N, and the parameters Le, N,  $\gamma$ ,  $\zeta$ ,  $F_0$ ,  $Ra_d$ ,  $\chi$ . One can conclude the following notes:



**Figure 19:** Variation of Nusselt number with Lewis number for various  $\gamma \square$  ( $\chi = 0.02$ . Ra<sub>d</sub> = 0.7, F<sub>0</sub> = 0.3, N = -0.1,  $\zeta = 0.0$ ).

With analogy to the case of absence of chemical reaction treated in [16], the non-Darcy parameter;  $F_0$  reduces the velocity close to the wall while it has no effect far from it, causes a slight increase in the temperature and reduces the heat transfer rate. Similarly, Lewis number; Le enhances the heat transfer rate, increases the velocity profile and decreases the concentration distribution. The buoyancy ratio; N increases the velocity, reduces both of temperature and concentration profiles and enhances both of heat and mass transfer rates.

It was found that the effect of chemical reaction is to decrease the concentration profile and increases the Sherwood number.

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# تأثير التفاعل الكيميائي والتشتت المزدوج على الانتقال الحراري والمادي في الأوساط المسامية اللادارسية

يعنى هذا البحث بدر اسة تأثير كل من التفاعل الكيميائي و التشتت الهيدروديناميكى المزدوج (الحراري والمادي) على الانتقال الحراري بالحمل الحر والانتقال المادي من سطح رأسي مطمور فى وسط مسامي لادارسى. ولقد أستخدم نموذج القوة الأسية للتعبير عن التفاعل الكيميائي فى معادلة التركيز. ولقد تم در اسة المسألة لكلا الوسط المسامي اللادارسى، والوسط المسامي الدارسى (يعتبر حالة خاصة من اللادارسى). ولقد تم إجراء تحويلات لا بعدية مناسبة لتحويل معادلات الحركة الحاكمة إلى مجموعة معادلات تفاضلية عادية. وتم حل هذه المعادلات مع تداخل مركب من ونسبة الطفو N، ومعامل التشتت الحراري الميكانيكي  $\gamma$ ، ومعامل التشت المادي الميكانيكي  $\zeta$ ، ومعامل التشتت الحراري الميكانيكي  $\gamma$ ، ومعامل التشت المادي ونسبة الطفو N، ومعامل التشتت الحراري الميكانيكي  $\gamma$ ، ومعامل التشت المادي الميكانيكي  $\zeta$ ، ومعامل التشت الحراري الميكانيكي والقد تم حساب السرعا الميكانيكي راحد معاد الادارسية التوليات الحراري الميكانيكي والقد تم والحرارية الفيزيائية الميكانيكي راحد والته معدل الانتقال الحراري الميكانيكي معادلات المادي ونسبة الطفو N، ومعامل التشت الحراري الميكانيكي والمال التشتال المادي الميكانيكي راحد معادلات تفاضلية اللابعدية الهامة مثل معدل الانتشار (عدد لويس المادي والميكانيكي راحيات الفيزيائية اللابعدية الهامة مثل معدل الانتشار (عدد لويا المادي والميكانيكي راحيات الفيزيائية اللابعدية الهامة مثل معدل الانتشار (عدد لويا المادي الميكانيكي راحيات الفيزيائية اللابعدية الهامة مثل معدل الانتشار (عدد لويا المادي والميكانيكي راحيات الفيزيائية اللابعدية الهامة مثل معدل الانتشان المادي