

The interaction between a Three-Level Atom and Four Systems of N-Two Level Atoms

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Abstract

We have four systems, each of them contains N-atoms, and each atom of them consists of two levels, and all of them interact with an atom consisting of three levels. Moreover, the atom is also related to the systems by coupling parameters that are time-dependent. By considering particular conditions, the exact solution for the wave function is realized. Then we survey the atomic population inversion in addition to the normalized correlation functions. We examine the influence of many parameters on the previous statistical aspects. We deduce from the conclusions that the behavior of the correlation function is affected by the existence of the time-dependent coupling parameters between spins. We prove that changing the values of the coupling parameters between spins λ_k , the parameters ε , and the quantum numbers $N\gamma$ enables us to control the correlated behavior in addition to the atomic population. So, we can control this model through these parameters.

keywords: Atom-Atom Interaction; Normalized Correlation Function; Non-Classical Behavior; Three-Level Atom.

1 Introduction

It is well known that the quantum mechanical atom-atom interaction can predict new aspects of the quantum nature of the system and the atom. There are a lot of models in quantum optics that describe the atom-atom interaction. For example, a single two-level atom interacts with N-two-level atoms with a magnetic field effect that has been discussed in Ref. [1-2] and other cases, such as Ref. [3-9]. Recently, much attention has focused on the influence of the time-dependent coupling parameters on various types of interactions. In the case of the atom-field interaction, the influence of time-dependent coupling parameters of the dimer system, has been discussed in Ref. [10]. Also, a three level Λ -type atom interacts with a two-mode with the existence of nonlinearity that has been discussed in Ref. [11- 12] and [13- 14]. On the other hand, in the atom-atom interaction, the effect of time-dependent coupling parameters of the model like our model, but for different cases from us that have been discussed recently in Ref. [15]. One of the most important goals of this paper is to survey the effect of atom-atom

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interaction on the atomic population inversion as well as the second-order correlation function. There are some examples of the atomic inversion evolution in the case of a two-level atom interacts with a various field, which has been discussed in Ref. [16 - 17]. Faraji et al. investigated the system of two-level atoms interacts with two fields by taking into account the dipole-dipole interaction of the atoms, and the time evolution of atomic inversion [18]. Faghihi et al., on the other hand, discussed the atomic inversion evolution of the system of a three-level atom of Λ -type interacts with a two-mode field in the existence of a bichromatic cavity [19]. Different types of interactions have been discussed in this field, for example, Ref. [20- 26]. Also, As the second-order function was evaluated, the system's non-classical behavior was better understood [27]- [32]. Dibakar et al. discussed many non-classical characteristics, like the sub-Poissonian statistics, correlation function among the two modes in Ref. [33]. Abdel-Wahab et al. examined the non-classical characteristic of a two two-level atoms that interacts with a two-mode cavity field [34]. There are different papers for several interactions and cases that have discussed the atomic inversion evolution and the second-order function as Ref. [35- 41].

In our previous paper [15], we studied the problem of four systems, each of them contains N -atoms, and each atom of them consists of two levels, and all of them interact with an atom consisting of three levels. Moreover, the atom is also related to the systems by coupling parameters that are time-dependent. We also investigated the influences of the quantum numbers N_γ ($\gamma = 1, 2, 3, 4$), the coupling parameters between spins λ_k ($k = 1, 2$), the parameter ε , and some initial cases on the previous statistical aspects.

In this paper, we explore the previous model, but for different Hamiltonian, and cases.

We arrange the paper in the following organization: In Section 2, we describe the model of the system and obtain the exact solution for the model describing the interaction between a three-level atom interacting with four systems of N -two level atoms with time-dependent coupling parameters, by taking some special cases. In section 3, and 4, we are investigating numerically the influences of the quantum numbers N_γ ($\gamma = 1, 2, 3, 4$), the coupling parameters between spins λ_k ($k = 1, 2$), the parameter ε , and the initial state on the previous statistical aspects. The conclusions are presented in section 5.

2 The Model

The following Hamiltonian explains our model as (we take $\hbar = 1$):

$$H = \sum_{\gamma=1}^4 \eta_{\gamma} J_z^{(\gamma)} + \sum_{l=1}^3 \Omega_l S_{ll} + \lambda_1(\tau) \left[S_{21} J_+^{(1)} J_+^{(2)} + J_-^{(2)} J_-^{(1)} S_{12} \right] + \lambda_2(\tau) \left[S_{32} J_+^{(3)} J_+^{(4)} + J_-^{(4)} J_-^{(3)} S_{23} \right], \quad (1)$$

Where $\eta_{\gamma}, \gamma = 1, 2, 3, 4$ are the frequencies of the systems, $\Omega_l, l = 1, 2, 3$ are the frequencies of the atom (where $\Omega_1 > \Omega_2 > \Omega_3$).

The atom coupling with the system by time-dependent coupling parameters $\lambda_k(\tau)$ ($k = 1, 2$) which given by:

$$\lambda_k(\tau) = \lambda_k \cos \varepsilon \tau, \quad k = 1, 2 \quad (2)$$

Also, $S_{\alpha\beta} = |\alpha\rangle\langle\beta|$ are the operators satisfy the commutation relation:

$$[S_{\alpha\beta}, S_{ij}] = S_{\alpha j} \delta_{i\beta} - S_{i\beta} \delta_{\alpha j} \quad (3)$$

The angular momentum $J_{\pm}^{(\gamma)}$ and $J_z^{(\gamma)}$ satisfy the relations:

$$J_L^{(\gamma)} = \frac{1}{2} \sum_{K_{\gamma}=1}^{N_{\gamma}} \sigma_L^{K_{\gamma}}, \quad L = x, y, z, \quad (4)$$

$$J_{\pm}^{(\gamma)} = J_x^{(\gamma)} \pm i J_y^{(\gamma)}, \quad (5)$$

$$[J_+^{(\gamma)}, J_-^{(l)}] = 2 J_z^{(\gamma)} \delta_{\gamma l} \quad (6)$$

$$[J_z^{(\gamma)}, J_{\pm}^{(l)}] = \pm J_{\pm}^{(\gamma)} \delta_{\gamma l} \quad (7)$$

And

$$J_{\pm}^{(\gamma)} |\zeta_{\gamma}, N_{\gamma}\rangle = \sqrt{(\zeta_{\gamma} \mp N_{\gamma})(\zeta_{\gamma} \pm N_{\gamma} + 1)} |\zeta_{\gamma}, N_{\gamma} \pm 1\rangle \quad (8)$$

$$J_{\pm}^{(\gamma)} |\zeta_{\gamma}, N_{\gamma}\rangle = N_{\gamma} |\zeta_{\gamma}, N_{\gamma}\rangle \quad (9)$$

We assume the state of our model to be in the form:

$$\begin{aligned} |W(\tau)\rangle &= G_1(\tau) |1, \zeta_1, N_1, \zeta_2, N_2, \zeta_3, N_3, \zeta_4, N_4\rangle \\ &+ G_2(\tau) |2, \zeta_1, N_1 + 1, \zeta_2, N_2 + 1, \zeta_3, N_3, \zeta_4, N_4\rangle \\ &+ G_3(\tau) |3, \zeta_1, N_1 + 1, \zeta_2, N_2 + 1, \zeta_3, N_3 + 1, \zeta_4, N_4 + 1\rangle \end{aligned} \quad (10)$$

The Schrodinger equation is given by:

$$i \frac{\partial}{\partial \tau} |W(\tau)\rangle = H |W(\tau)\rangle. \quad (11)$$

From equations (1), (10), and (11), we have the following differential equations:

$$\begin{aligned}
 i \frac{dG_1(\tau)}{d\tau} &= \beta_1 G_1(\tau) + \Gamma_1 \cos \varepsilon \tau G_2(\tau), \\
 i \frac{dG_2(\tau)}{d\tau} &= \beta_2 G_2(\tau) + \Gamma_1 \cos \varepsilon \tau G_1(\tau) + \Gamma_2 \cos \varepsilon \tau G_3(\tau), \\
 i \frac{dG_3(\tau)}{d\tau} &= \beta_3 G_3(\tau) + \Gamma_2 \cos \varepsilon \tau G_2(\tau),
 \end{aligned} \tag{12}$$

Where,

$$\beta_1 = \eta_1 N_1 + \eta_2 N_2 + \eta_3 N_3 + \eta_4 N_4 + \Omega_1, \tag{13}$$

$$\beta_2 = \eta_1(N_1 + 1) + \eta_2(N_2 + 1) + \eta_3 N_3 + \eta_4 N_4 + \Omega_2, \tag{14}$$

$$\beta_3 = \eta_1(N_1 + 1) + \eta_2(N_2 + 1) + \eta_3(N_3 + 1) + \eta_4(N_4 + 1) + \Omega_3, \tag{15}$$

$$\Gamma_1 = \lambda_1 \sqrt{(\zeta_1 - N_1)(\zeta_1 + N_1 + 1)(\zeta_2 - N_2)(\zeta_2 + N_2 + 1)} \tag{16}$$

$$\Gamma_2 = \lambda_2 \sqrt{(\zeta_3 - N_3)(\zeta_3 + N_3 + 1)(\zeta_4 - N_4)(\zeta_4 + N_4 + 1)} \tag{17}$$

Now, we can solve the system (12), by these steps:

$$\dot{G}(\tau) = F(\tau)G(\tau) \tag{18}$$

Where

$$G(\tau) = \begin{pmatrix} G_1 \\ G_2 \\ G_3 \end{pmatrix}, F(\tau) = \begin{pmatrix} -i\beta_1 & -i\Gamma_1 \cos \varepsilon \tau & 0 \\ -i\Gamma_1 \cos \varepsilon \tau & -i\beta_2 & -i\Gamma_2 \cos \varepsilon \tau \\ 0 & -i\Gamma_2 \cos \varepsilon \tau & -i\beta_3 \end{pmatrix} \tag{19}$$

By considering particular conditions, the matrix $F(\tau)$ of the system is symmetric and the product of the matrix and the integral of this matrix is commutative. Therefore, the fundamental matrix of the system is given by:

$$Q(\tau) = \exp \left(\int_0^\tau F(\tau) d\tau \right) \tag{20}$$

We evaluate the matrix exponential by converting the matrix to diagonal form, also we compute the eigenvalues and eigenvectors of the matrix $(\int_0^\tau F(\tau)d\tau)$. Then, we get the transition matrix (M) of the eigenvectors. The Jordan form R of the matrix can be written as:

$$R = M^{-1} \left[\int_0^\tau F(\tau)d\tau \right] M \quad (21)$$

Therefore, we can compute the fundamental matrix as:

$$Q(\tau) = M^{-1} \exp(R) M \quad (22)$$

Then the solution become $G(\tau) = Q(\tau)C$, where $C = G(0)$. From the initial conditions, the unknown vector $G(\tau)$ can be obtained, then the nonhomogeneous system (12) can be solved.

Then, we get the solution of the system as follows:

$$G_1(\tau) = \frac{e^{-i\beta_1\tau}}{\nu^2} \{(\Gamma_1^2 \cos \theta + \Gamma_2^2)G_1(0) - i\nu\Gamma_1 \sin \theta G_2(0) + \Gamma_1\Gamma_2(\cos \theta - 1)G_3(0)\}, \quad (23)$$

$$G_2(\tau) = -\frac{e^{-i\beta_1\tau}}{\nu} \{i \sin \theta (\Gamma_1 G_1(0) + \Gamma_2 G_3(0)) - \nu \cos \theta G_2(0)\}, \quad (24)$$

$$G_3(\tau) = \frac{e^{-i\beta_1\tau}}{\nu^2} \{(\cos \theta - 1)\Gamma_1\Gamma_2 G_1(0) - i\nu\Gamma_2 \sin \theta G_2(0) + (\Gamma_2^2 \cos \theta + \Gamma_1^2)G_3(0)\}, \quad (25)$$

Where

$$\theta = \nu\epsilon(\tau), \quad (26)$$

$$\epsilon(\tau) = \frac{\sin(\epsilon\tau)}{\epsilon}, \quad (27)$$

$$\nu = \sqrt{\Gamma_1^2 + \Gamma_2^2}. \quad (28)$$

3 The Atomic Inversion

The atomic inversion can be written as follows [42]- [45]:

$$\langle S_z \rangle = |G_1(\tau)|^2 + |G_2(\tau)|^2 - |G_3(\tau)|^2, \quad (29)$$

To discuss the atomic inversion, we have plotted several figures in which we take into account $\zeta_1 = 50$, $\zeta_2 = 40$, $\zeta_3 = 30$, $\zeta_4 = 20$, and various values of λ_k , N_γ , ε , also at initially the state be {i} $|W(0)\rangle_r = |1\rangle$, {ii} $|W(0)\rangle_r = \frac{1}{\sqrt{2}}|1\rangle + \frac{1}{\sqrt{2}}|2\rangle$, and {iii} $|W(0)\rangle_r = \frac{1}{\sqrt{2}}|1\rangle + \frac{1}{2}|2\rangle + \frac{1}{2}|3\rangle$. We show that in all figures the oscillations of the atomic inversion curve are periodic. In **Fig. 1**, we modify the values of the coupling parameters λ_k and make the other parameters fixed i.e. $N_\gamma = 1, \varepsilon = 1$, also at initially state be {i} to examine their influences on the atomic inversion. We note that the number of oscillations rises and the maximum value, $\langle S_z \rangle = 1$, reductions to attain its minimum value, $\langle S_z \rangle = -0.2$, as the coupling parameters λ_k rise. We have at the top of the curve constant interval at $\lambda = 0.25$ decay as the parameters λ_k rise. In **Figs. 2(a-c)**, we modify the values of the quantum numbers N_γ and make the other parameters fixed i.e. $\lambda_k = 1, \varepsilon = 1$, to examine their influences on the atomic inversion. As the quantum numbers, N_γ rise the number and the phase of oscillations reduction. In **Fig. 2(a)**, initially state be {i}, but in **Fig. 2(b)**, initially state be {ii}. In both, the curve begins at the highest point, $\langle S_z \rangle = 1$, this high point reductions to attain the lowest point, $\langle S_z \rangle = 0.3$, in **Fig. 2(a)** at $N_\gamma = 1$. On the other hand, in **Fig. 2(b)**, the high point reductions to attain the lowest point but at $\langle S_z \rangle = 0.7$, at $N_\gamma = 1$. As the quantum numbers N_γ rise, the lowest point rises to attain $\langle S_z \rangle \simeq 0.9$ at $N_\gamma = 19$, in **Fig. 2(a)**. Also in **Fig. 2(b)**, as the quantum numbers N_γ rise, the lowest point rises to be closed too 1. But in **Fig. 2(c)**, initially state be {iii}, the curve begins at the lowest point, $\langle S_z \rangle = 0.5$, this lowe point rises to attain the highest point, $\langle S_z \rangle = 1$, at $N_\gamma = 1$. As the quantum numbers N_γ rise, the maximum value reduction to attain $\langle S_z \rangle = 0.7$ at $N_\gamma = 19$. in **Figs. 3(a-c)**, we modify the value of the parameter ε and make the other parameters fixed i.e. $\lambda_k = 1, N_\gamma = 1$, to examine their influences on the atomic inversion. As the parameter ε rises the values rise, but phase of oscillations reductions. in **Fig. 3(a)**, initially state be {i}, the curve begins at the highest point, $\langle S_z \rangle = 1$, this high point reductions to attain the lowest point, $\langle S_z \rangle \simeq 0.35$. As the parameter ε rises, the lowest point rises to attain $\langle S_z \rangle \simeq 0.65$. Also in **Fig. 3(b)**, initially state be {ii}, the curve begins at the highest point, $\langle S_z \rangle = 1$, this the high point reductions to attain the lowest point, $\langle S_z \rangle \simeq 0.65$. As the parameter ε rises, the lowest point rises to attain $\langle S_z \rangle \simeq 0.75$. But in

Fig. 3(c), initially state be {iii} the curve begins at the lowest point, $\langle S_z \rangle = 0.5$, this lowest point rises to attain the highest point, $\langle S_z \rangle = 1$. As the parameter ε rises, the highest point reduction to attain $\langle S_z \rangle \approx 0.95$.

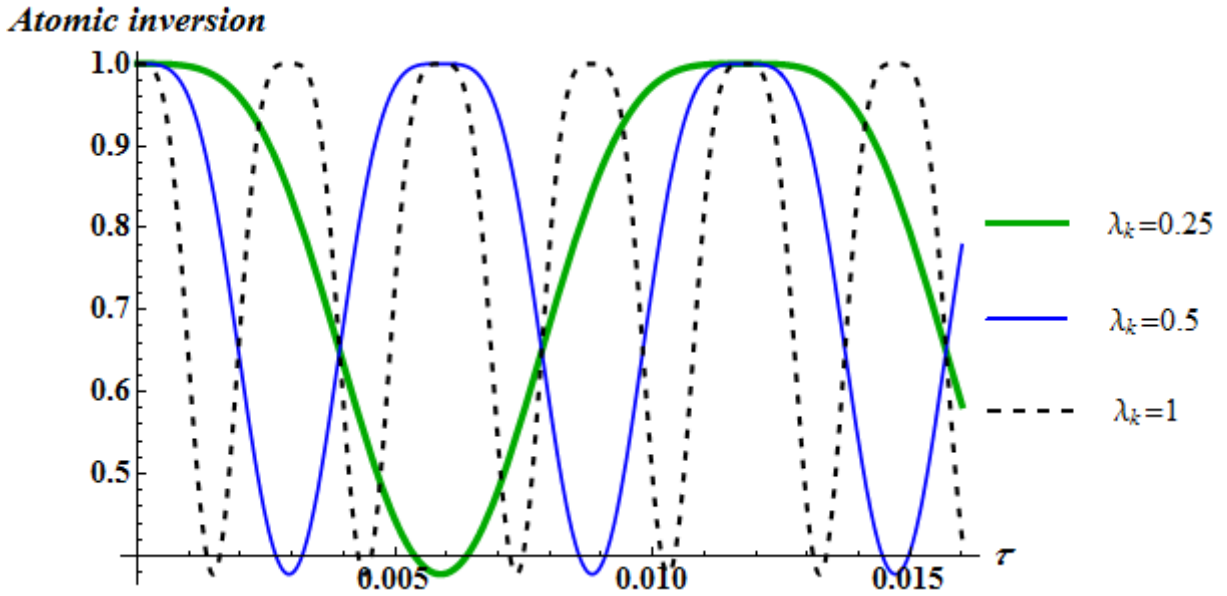
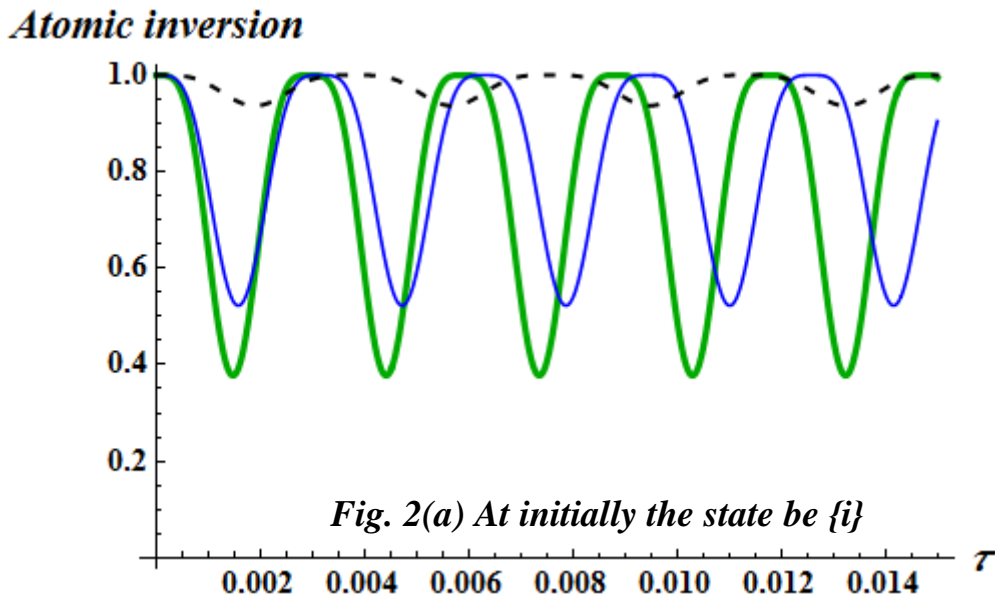


Fig. 1: The atomic inversion evolution at initially the state be {i} with $N_\gamma = 1$, and $\varepsilon = 1$.



Atomic inversion

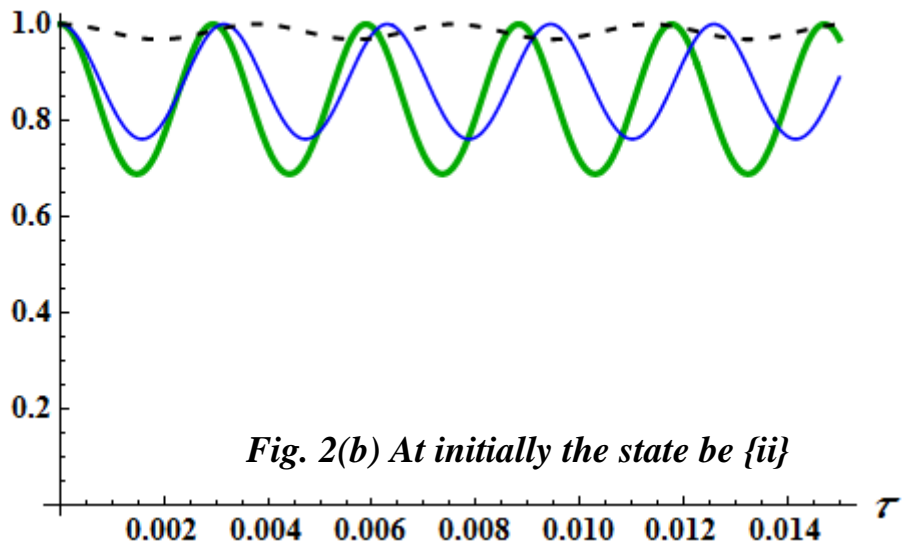


Fig. 2(b) At initially the state be {ii}

Atomic inversion

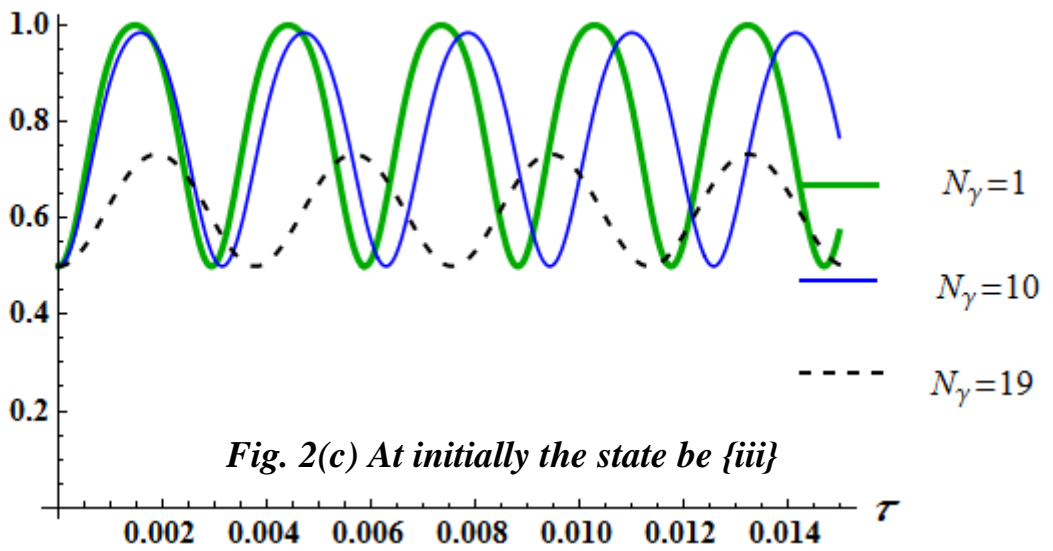
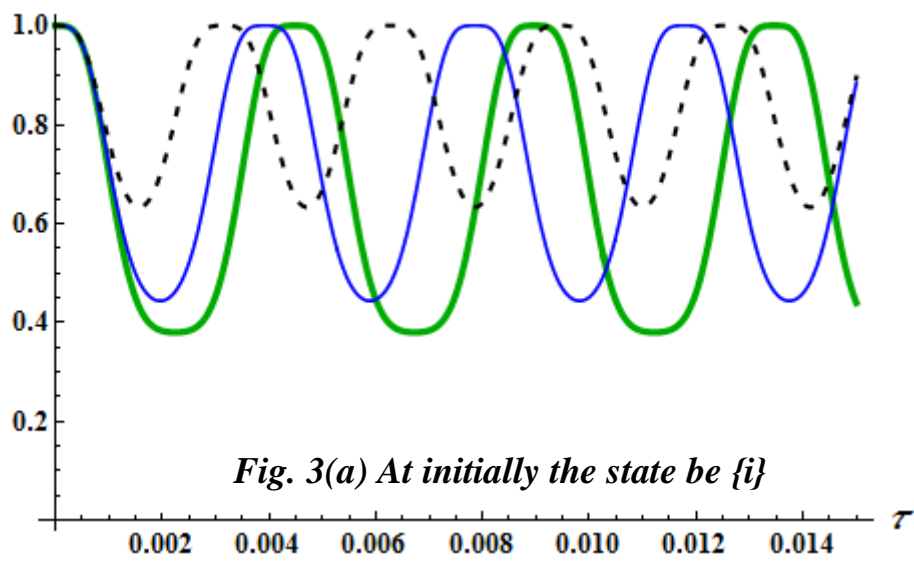


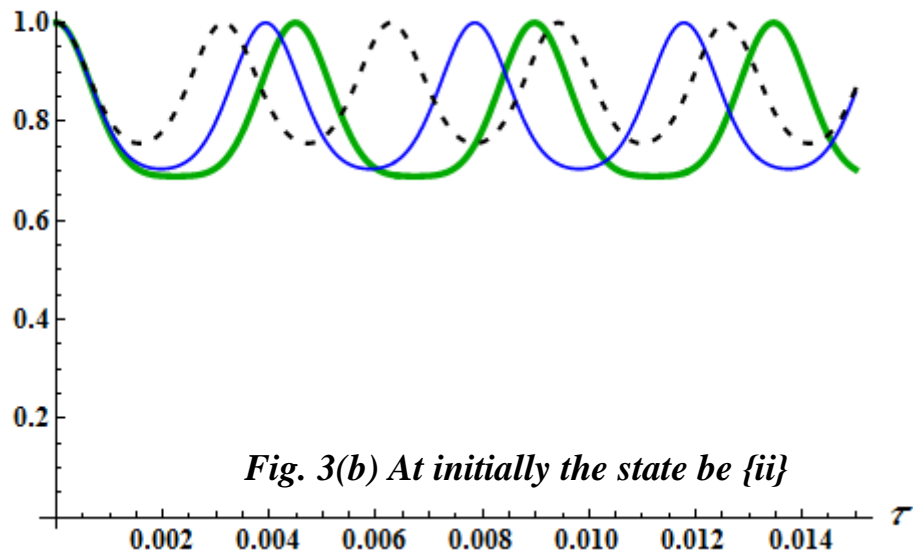
Fig. 2(c) At initially the state be {iii}

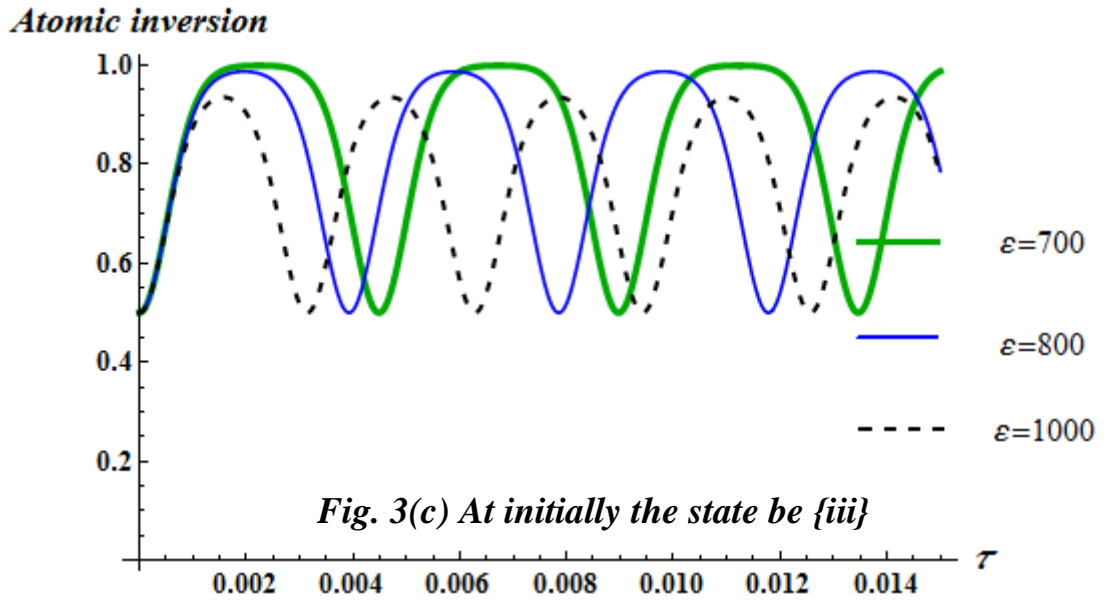
Figs. 2: The figures for $\lambda_k = 1$, and $\varepsilon = 1$.

Atomic inversion



Atomic inversion





Figs. 3: The figures for $\lambda_k = 1$, and $N_\gamma = 1$.

4 The Correlation Function

The normalized second-order correlation function can be written as [2], [15], [24] and [46]:

$$g_{\beta}^2(\tau) = \frac{\langle J_+^{(\beta)2} J_-^{(\beta)2} \rangle}{\langle J_+^{(\beta)} J_-^{(\beta)} \rangle^2}, \quad \beta = 1, 2, 3, 4. \quad (30)$$

Then we evaluate these expectation values, where

$$J_+^{(\beta)2} J_-^{(\beta)2} = \varphi(J_z^{(\beta)} - 1) \varphi(J_z^{(\beta)}), \quad (31)$$

$$J_+^{(\beta)} J_-^{(\beta)} = \varphi(J_z^{(\beta)}), \quad (32)$$

Equation (31) becomes

$$g_{\beta}^2(\tau) = \frac{\langle \varphi(J_z^{(\beta)} - 1) \varphi(J_z^{(\beta)}) \rangle}{\langle \varphi(J_z^{(\beta)}) \rangle^2}, \quad (33)$$

$$\varphi(J_z^{(\beta)}) = J^{(\beta)2} - J_z^{(\beta)2} + J_z^{(\beta)2}. \quad (34)$$

By substituting from (10) in (31) and (32), we have:

$$g_1^2(\tau) = \frac{\varphi(N_1, \zeta_1) \{ \varphi(N_1 - 1, \zeta_1) |G_1(\tau)|^2 + \varphi(N_1 + 1, \zeta_1) [|G_2(\tau)|^2 + |G_3(\tau)|^2] \}}{\{ \varphi(N_1, \zeta_1) |G_1(\tau)|^2 + \varphi(N_1 + 1, \zeta_1) [|G_2(\tau)|^2 + |G_3(\tau)|^2] \}^2}, \quad (35)$$

$$g_2^2(\tau) = \frac{\varphi(N_2, \zeta_2)\{\varphi(N_2 - 1, \zeta_2)|G_1(\tau)|^2 + \varphi(N_2 + 1, \zeta_2)[|G_2(\tau)|^2 + |G_3(\tau)|^2]\}}{\{\varphi(N_2, \zeta_2)|G_1(\tau)|^2 + \varphi(N_2 + 1, \zeta_2)[|G_2(\tau)|^2 + |G_3(\tau)|^2]\}^2}, \quad (36)$$

$$g_3^2(\tau) = \frac{\varphi(N_3, \zeta_3)\{\varphi(N_3 - 1, \zeta_3)[|G_1(\tau)|^2 + |G_2(\tau)|^2] + \varphi(N_3 + 1, \zeta_3)|G_3(\tau)|^2\}}{\{\varphi(N_3, \zeta_3)[|G_1(\tau)|^2 + |G_2(\tau)|^2] + \varphi(N_3 + 1, \zeta_3)|G_3(\tau)|^2\}^2}, \quad (37)$$

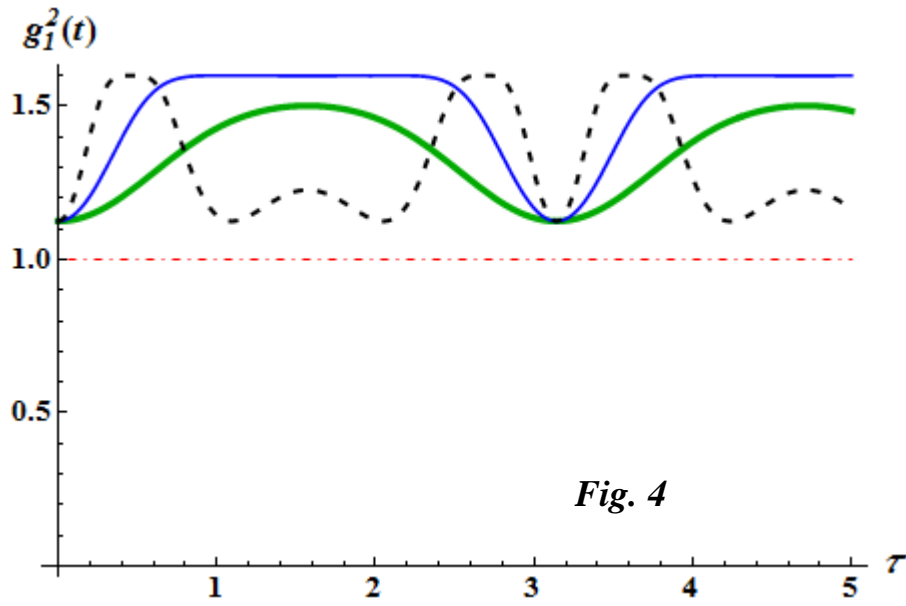
$$g_4^2(\tau) = \frac{\varphi(N_4, \zeta_4)\{\varphi(N_4 - 1, \zeta_4)[|G_1(\tau)|^2 + |G_2(\tau)|^2] + \varphi(N_4 + 1, \zeta_4)|G_3(\tau)|^2\}}{\{\varphi(N_4, \zeta_4)[|G_1(\tau)|^2 + |G_2(\tau)|^2] + \varphi(N_4 + 1, \zeta_4)|G_3(\tau)|^2\}^2}, \quad (38)$$

Where

$$\varphi(N_\beta, \zeta_\beta) = (\zeta_\beta + N_\beta)(\zeta_\beta - N_\beta + 1), \quad \beta = 1,2,3,4 \quad (39)$$

We note that from pervious equations that $g_1^2(\tau) = g_2^2(\tau)$ also $g_3^2(\tau) = g_4^2(\tau)$. So, Figs. 4,5 and 6 show the normalized second-order correlation functions $g_1^2(\tau)$, and $g_3^2(\tau)$ only vs. the time τ . To discuss the normalized second-order correlation functions we will plot several figures in which we take into account $\zeta_\gamma = \frac{5}{2}$, also at initially state be $\{i\}$ and modify different values of λ_k , N_γ , and ε parameters. We show that in all figures the oscillations of the second-order correlation functions is periodic and regular. In **Fig. 4**, we modify the values of the coupling parameters λ_k and make the other parameters fixed i.e. $N_\gamma = 1$ and $\varepsilon = 1$, to examine their influences on the correlation functions $g_1^2(\tau)$ and $g_3^2(\tau)$. The correlation functions $g_1^2(\tau)$, $g_3^2(\tau) > 1$, thus these functions show super-Poissonian behavior. In **Fig. 5**, we also modify the values of the quantum numbers N_γ and make the other parameters fixed i.e. $\lambda_k = 1$ and $\varepsilon = 1$, to examine their influences on the correlation functions $g_1^2(\tau)$ and $g_3^2(\tau)$. The correlation functions $g_1^2(\tau)$ and $g_3^2(\tau)$ grow up As N_γ rises. At $N_\gamma = \frac{-3}{2}, \frac{-1}{2}$ the functions $g_1^2(\tau)$ and $g_3^2(\tau) < 1$, so these functions show sub-Poissonian behavior. But As $N_\gamma = \frac{1}{2}$, the correlation functions $g_1^2(\tau)$ and $g_3^2(\tau)$ sometimes show sub-Poissonian behavior ($g_1^2(\tau), g_3^2(\tau) < 1$) and sometimes show super-Poissonian behavior ($g_1^2(\tau), g_3^2(\tau) > 1$). As $N_\gamma = \frac{3}{2}$, the correlation functions $g_1^2(\tau)$ and $g_3^2(\tau)$ show super-Poissonian behavior ($g_1^2(\tau), g_3^2(\tau) > 1$). In **Fig. 6**, we modify the values of the parameter ε and make the other parameters fixed i.e. $N_\gamma = 1$ and $\lambda_k = 1$, to examine their influences on the correlation functions $g_1^2(\tau)$ and $g_3^2(\tau)$. The correlation functions $g_1^2(\tau), g_3^2(\tau) > 1$ at all values of the parameter ε ,

thus these functions show super-Poissonian behavior. We note that the number of oscillations rises as the parameter ε increases, yet the phase of the oscillations reductions.



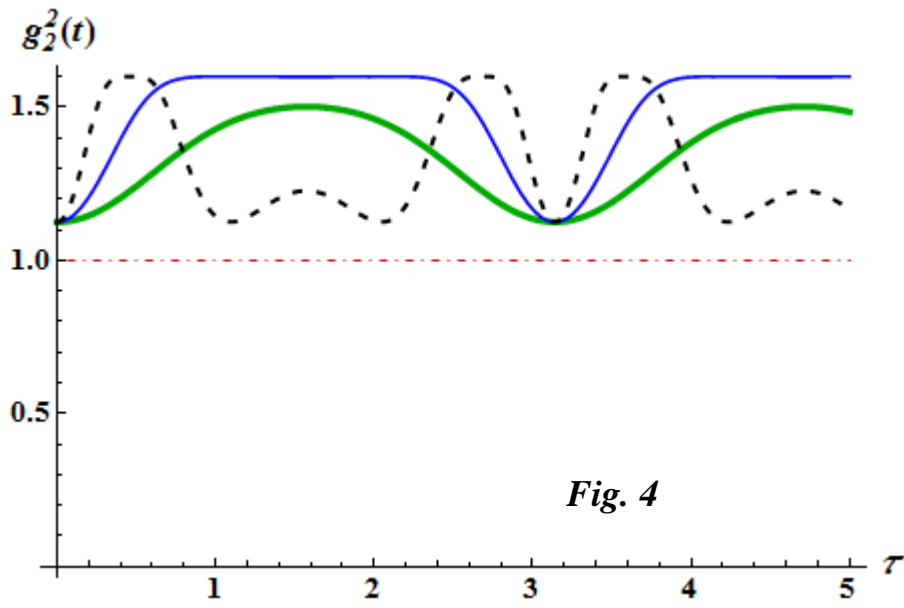


Fig. 4

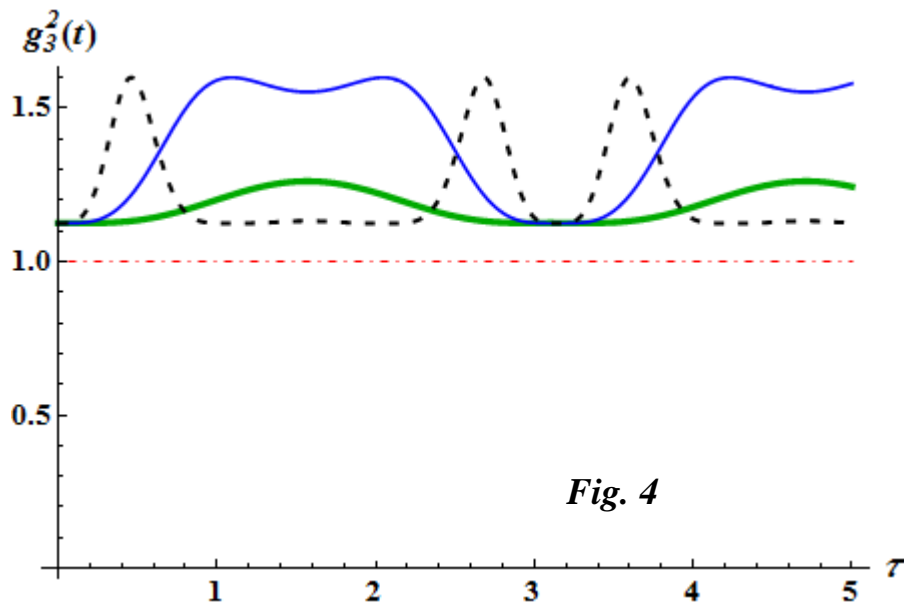
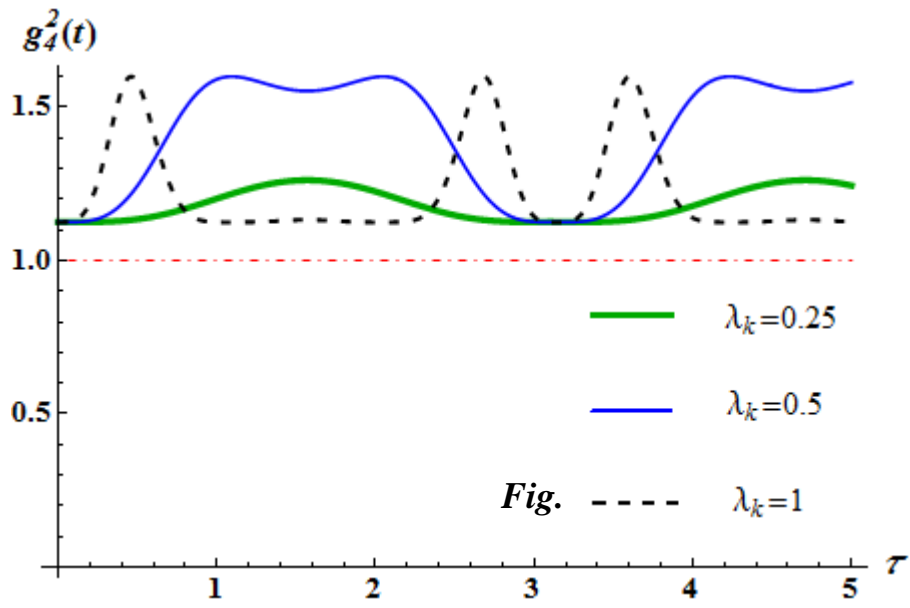
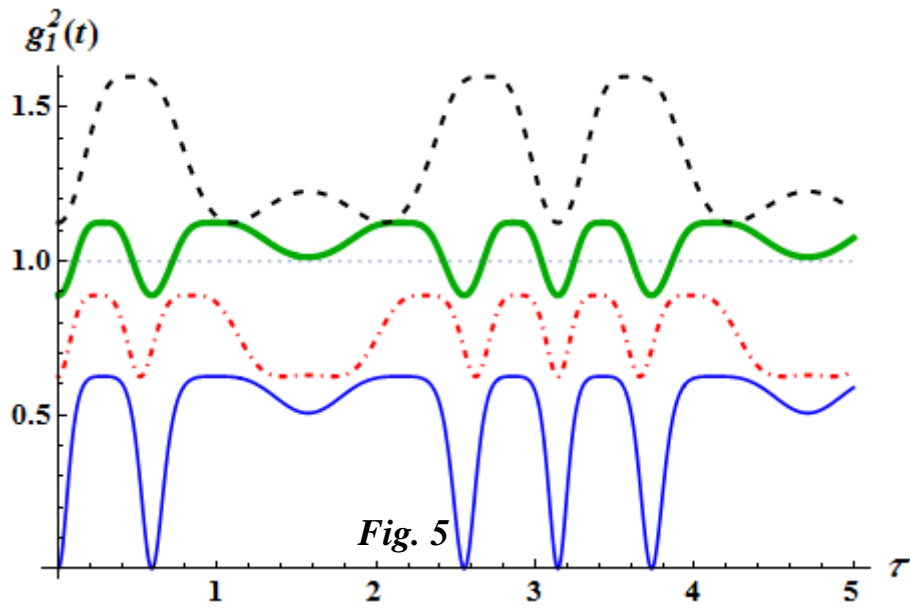
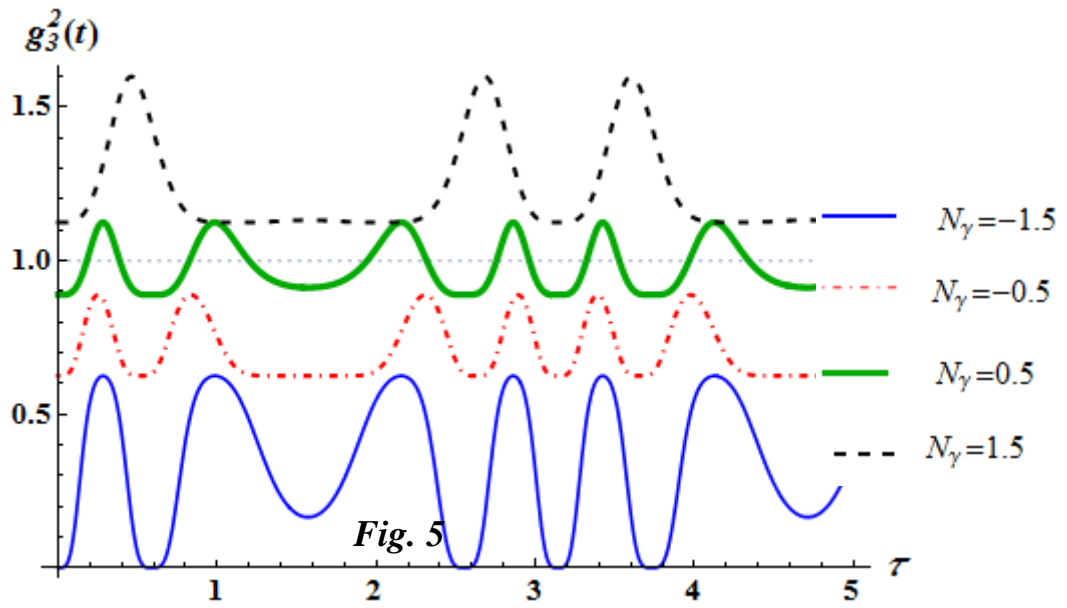


Fig. 4

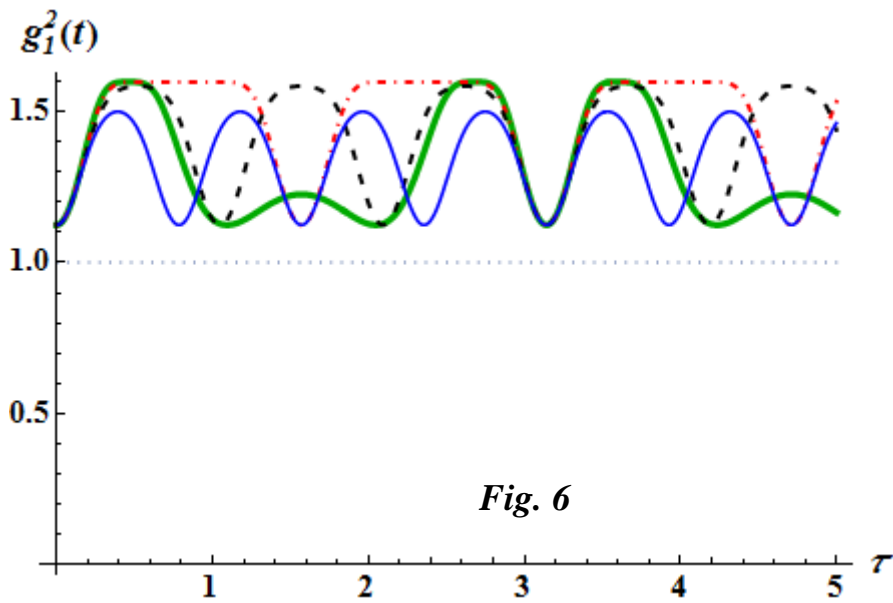


Figs. 4: The functions $g_{\beta}^2(\tau)$ vs. the time τ . At initially the state be $\{i\}$ with $N_{\gamma} = \frac{3}{2}, \varepsilon = 1$





Figs. 5: The functions $g_{\beta}^2(\tau)$ vs. the time τ . At initially the state be $\{i\}$ with $\lambda_K = 1, \varepsilon = 1$.



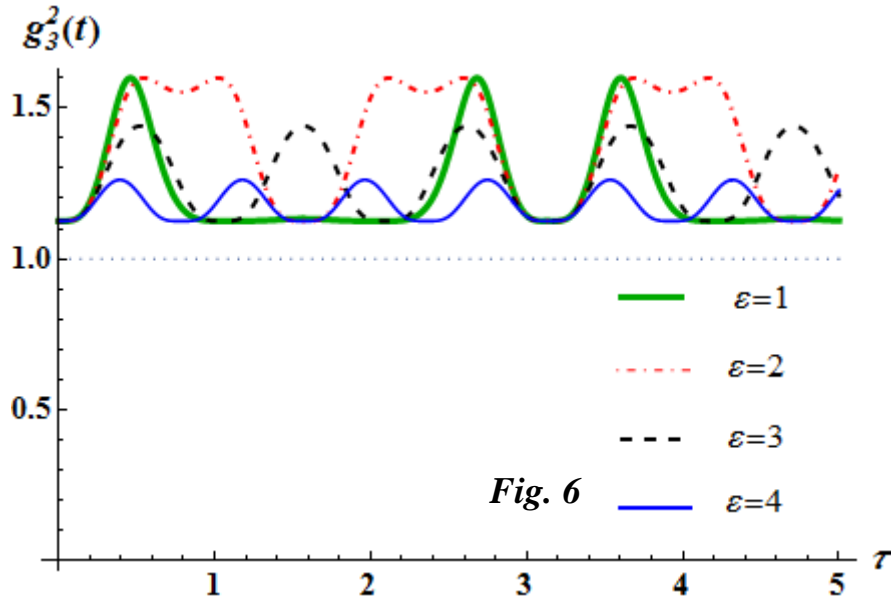


Fig. 6

Figs. 6: The functions $g_{\beta}^2(\tau)$ vs. the time τ . At initially the state be $\{i\}$ with $N_{\gamma} = \frac{3}{2}, \lambda_K = 1$.

5 Conclusion

In this paper, we have studied four systems, each of them contains N-atoms, and each atom of them consists of two levels, and all of them interact with an atom consisting of three levels. The coupling parameters between the atom and the systems are modulated to be time-dependent. By considering particular conditions, the exact solution for this model has been obtained. Therefore, the wave function has been obtained as well. We also have discussed the influences of various parameters like for example the coupling parameters λ_k , the quantum numbers N_{γ} , and the (decay rates) parameter ε on some statistical aspects. We have managed to study the atomic population inversion in addition to the second-order correlation function. In our handling of the atomic population inversion, we have realized that there is a constant interval in the top of the curve at $\lambda = 0.25$, and then it decays and the maximum value, $\langle S_z \rangle = 1$, reductions to attain its minimum value, $\langle S_z \rangle = -0.2$, as the coupling parameters λ_k rise. On the other hand, in Ref. [2-15], this constant interval also decays as the coupling parameters λ rise and the number of periodic oscillations rises, also the maximum value, $\langle S_z \rangle = 1$. When the coupling parameters λ rise, the maximum value reductions to attain a different minimum value than our model, $\langle S_z \rangle = -0.75$, in Ref.[2], and, $\langle S_z \rangle = -0.4$, in Ref.[15]. This difference in the minimum value between the three cases can be assigned to a different number of systems and the Hamiltonian. We show that as the parameter ε rises for all initial be $\{i\}$, $\{ii\}$, and $\{iii\}$, the

values rise, but phase of oscillations reductions. For the first two initial states, the minimum value of the atomic population inversion curve rises to attain $\langle S_z \rangle \simeq 0.65$. But at the last initial state, the maximum value of the atomic population inversion curve reductions to attain $\langle S_z \rangle \simeq 0.95$. Also, we note that the correlation function $g_1^2(\tau) = g_2^2(\tau)$ and also $g_3^2(\tau) = g_4^2(\tau)$. The correlation functions at all values of the parameter ε and the coupling parameters λ_k , thus these functions show super-Poissonian behavior. We note that the number of oscillations rises as the parameter ε rises, but the phase of the oscillations reductions. On the other hand, As changing the quantum numbers N_γ , the correlation functions $g_1^2(\tau)$ and $g_3^2(\tau)$ sometimes show sub-Poissonian behavior ($g_1^2(\tau), g_3^2(\tau) < 1$) and sometimes show super-Poissonian behavior ($g_1^2(\tau), g_3^2(\tau) > 1$). The correlation functions $g_1^2(\tau)$ and $g_3^2(\tau)$ grow up As N_γ rises. Finally, we have proved that changing the values of the coupling parameters between spins λ_k , the parameters ε , and the quantum numbers N_γ enables us to control the correlated behavior.

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التفاعل بين ذرة ثلاثية المستويات وأربعة أنظمة تحتوي على ن من الذرات ثنائي المستوي

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نبذة مختصرة

لدينا أربعة أنظمة، كل منها يحتوي على ن من الذرات، وكل ذرة منها تتكون من مستويين، وكلها تتفاعل مع ذرة تتكون من ثلاثة مستويات. علاوة على ذلك، ترتبط الذرة أيضًا بالأنظمة من خلال معاملات الاقتران التي تعتمد على الزمن. من خلال حالات خاصة، يتم الحصول على الحل الدقيق للدالة الموجة. ثم نقوم بدراسة الانقلاب الذري بالإضافة إلى الارتباط الطبيعي. ندرس تأثير العديد من المعاملات على الجوانب الإحصائية السابقة. نستنتج من ذلك أن سلوك دالة الارتباط يتأثر بوجود معاملات الاقتران التي تعتمد على الزمن بين الدورات. لذلك، يمكننا التحكم في هذا النموذج من خلال هذه المعاملات.