Electrical and Optical Properties of Ternary System of Bi-Sb-Se

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Polycrystalline alloys of $Bi_{48-x}Sb_xSe_{52}$ (where x = 0.6, 11, 25, 35 and 44 at.%) have been prepared by direct fusion technique. The crystalline nature has been identified using X-ray diffraction. The electrical measurements of dc and ac conductivities as well as the infrared transmittance were investigated. Compositional dependences of electrical conductivity and the absorption edge are discussed from the stand point of chemical bonds formed in alloys and related to changes of the free charge carriers concentration (Moss – Burstein effect) and ionicity due to incorporation of Sb atoms in high concentrations. The data of the UV/VIS spectroscopy showed that the mechanism of the optical absorption followed the rule of allowed direct transition for all compositions.

1. Introduction:

Ternary compound semiconductors are currently being investigated due to their variety of applications in thermoelectric devices at room temperature, photovoltaic solar cell materials, electro- optics, opto-electronics and non-linear optics. Solid solutions based on bismuth and antimony chalcogenides find a wide range of thermoelectric applications now as high-efficiency devices [1-3]. Candidate materials for thermoelectric applications should possess high electrical conductivity, high Seebeck coefficient, and low thermal conductivity. Considerable attention have been devoted to the research of Bi₂Se₃ [4,5]. Some studies on optical and electrical conductivity measurements of the single crystals (Bi_{1-x}Sb_x)₂Se₃ for x = 0 to 0.2 have been reported [6,7]. In this article the compositional dependence of the electrical conductivity and the optical energy gap of Bi_{48-x}Sb_xSe₅₂ (where x = 0.6, 11, 25, 35 and 44 at.%) polycrystalline alloys is discussed.

2. Experimental:

Polycrystalline alloys of Bi_{48-x}Sb_xSe₅₂ (where x = 0.6, 11, 25, 35 and 44 at.%) have been prepared by direct fusion technique. Proper amounts of Sb, Bi and Se (purity 99.999 %) were mixed together and put in silica tubes which are then sealed off at vacuum $\approx 10^{-5}$ Torr. The sealed ampoules were inserted in an electric furnace held at 1173 K for 24h. During the course of heating, the tubes were shaken vigorously at regular intervals to ensure homogeneity of the alloys. The tubes were then cooled gradually to room temperature. Five specimens of as-prepared samples were ground up and put onto ultrasonically cleaned glass substrates to be examined using X-ray diffractometer, type,[X" PERT-MPDUG (Philips PW 1370)] to confirm the morphology and crystalline nature of materials. The results of energy dispersive X-ray (EDX) confirmed the presence of Sb, Bi and Se in the proper amounts and no other impurities exist.

The infrared data was recorded at normal incidence in the wavelength range 2.5-25 μ m by means of Fourier Transform Infrared Spectrometer [Jasco FT/IR- 410]. The data of transmittance (T) and absorbance (A) were determined at normal incidence in the wavelength range 0.3-1.1 μ m by means of double-beam UV/VIS spectrophotometer (SPECORD 40/analytik-jena).

An electrometer (Keithly 6517A) was used for measuring the electrical conductivity by measuring the electrical resistance of each specimen which fabricated in the disk form of 1 cm diameter and 3 mm thickness. The specimen was held between two graphite electrodes (as pressure contact to it) in a copper holder. The sample holder was placed in an electric furnace in which the temperature was varied by heating rates 1 K/min using a temperature controller. The dielectric properties were measured using a bridge (type INSTEK-LCR 821) in the frequency range 10^2 -2 x 10^5 Hz. The cross-linked samples were in the form of disks of range 0.156 to 0.262 cm thick. The samples were sandwiched between two brass electrodes and both their capacitance and their dielectric loss tangent (tan δ) values were measured at different frequencies. The dielectric constant ϵ ` (real part of the dielectric constant) of the samples and the a.c conductivity σ_{ac} were calculated by using the relations:

$$\varepsilon' = \frac{d}{\varepsilon \cdot A} C \tag{1}$$

$$\sigma_{ac} = \omega \varepsilon_{\circ} \varepsilon'' \tag{2}$$

$$\mathbf{\epsilon}^{\prime\prime} = \mathbf{\epsilon}^{\prime} \tan \mathbf{\delta} \tag{3}$$

where C is the capacitance, d is the thickness of the sample and A is the crosssectional area of each of the parallel surface of the sample. ε_{\circ} is the permittivity of free space (8.85x10⁻¹² F\m), ε'' is the dielectric loss (imaginary part of the dielectric constant) and ω is the angular frequency ($\omega = 2\pi f$).

3. Results and Discussion: 3.1. X-ray Diffraction:

X-ray diffraction patterns of as-deposited $Bi_{48-x}Sb_xSe_{52}$ (where x = 0.6, 11, 25, 35 and 44 at. %) alloys were recorded at room temperature. The crystalline phases of the alloys were detected from the corresponding peaks appeared in the X-ray patterns as shown in Fig. (1) which indicate that the alloys are polycrystalline in nature. The analysis of the diffraction pattern according to the ASM cards showed that the phases appeared are rhombohedral (Bi₂Se₃, Bi and Sb), orthorhombic (Sb₂Se₃) and monoclinic (Se). It is seen that with increasing the Sb content, the peak intensity corresponding to the Sb₂Se₃ and Sb phases increases while that corresponding to Bi₂Se₃ phases decreases sharply. This indicate that the Sb₂Se₃ phases grows in the matrix at the expense of Bi₂Se₃ phases as the Sb content increases.



Fig. (1): X-ray diffraction patterns of Bi_{48-x}Sb_xSe₅₂ alloys.

3.2. Electrical Properties: 3.2.1. DC Conductivity:

Figure (2) show the temperature dependence of the d.c conductivity for the $Bi_{48-x}Sb_xSe_{52}$ (where x = 0.6, 11, 25, 35 and 44 at. %), in the temperature range 303 – 393 K. All the curves show straight lines for the plots of (ln σ) vs. (1/T) obeying the relation:

$$\sigma_{\rm d} = \sigma_{\rm o} \, {\rm e}^{-{\rm E}_{\rm g}/{\rm k}{\rm T}} \tag{4}$$

where E_a is the activation energy, k is the Boltzmann's constant, while σ_o is constant.

The results exhibit only one type of conduction mechanism in the temperature range 303 - 393 K for all compositions. This mechanism is regarded as a regular band – type conduction in the localized states. The activation energy of conduction and dc conductivity measured at 303 K are summarized in table 1 for all compositions. It is seen that the dark conductivity increases up to x = 11 at. % and then decreases with increasing the Sb content, while the activation energy changes with Sb content in an opposite manner.

Sample No.	Composition(x) at.%	E _a (eV)	$\sigma_{303K}(\Omega^{-1}cm^{-1})$
A4	0.6	0.085	0.04
A12	11	0.029	0.12
A20	25	0.188	0.004
A28	35	0.478	1.33E-6
A36	44	0.523	7.80E-5

Table (1): The values of δ_{303K} and E_a for $Bi_{48-x}Sb_xSe_{52}$ alloys.

For single crystals of $(Bi_{1-x}Sb_x)_2Se_3$ solid solutions (x = 0 - 0.2) [7], it has been reported that the incorporation of the antimony atoms in Bi_2Se_3 crystal lattice results in an increase of the free electron concentration for low antimony content, whereas the free electron concentration is suppressed in the range of high antimony content [7]. In agreement, for the system $Bi_{48-x}Sb_xSe_{52}$ under study, the measurements showed an increase in conductivity for x = 0.6 and 11 at. % and a decrease in conductivity for higher antimony concentration. This means that the free electron concentration decreases with rising the Sb content from 25 up to x = 44 at. %.



Fig. (2): Temperature dependence of the dark conductivity σ_{dc} for the $Bi_{48-x}Sb_xSe_{52}$ bulk alloys.

3. 2.2. A.C conductivity Measurements:

The plot of a.c. conductivity as a function of frequency for the system $Bi_{48-x}Sb_xSe_{52}$, where x = 0.6, 25, 35 and 44 at.% is illustrated in Fig. (3). It is seen that the a.c. conductivity is decreased with rising the frequency for all compositions, except for x = 44 at. %, it seems to be constant.



Fig. (3): AC conductivity vs. the frequency for Bi_{48-x}Sb_xSe₅₂ alloys.

in the range of high antimony content.

The frequency dependence of the dielectric constant ε' and dielectric loss ε'' are shown in Figs. (4 and 5), respectively. It can be observed that all curves obtained for the system Bi_{48-x}Sb_xSe₅₂, where x = 0.6, 25, 35 and 44 at.% show that ε ` and ε `` decrease almost linearly with increasing the frequency in the range $10^3 - 10^5$ Hz and remain almost constant at higher frequencies. One can conclude that the real and imaginary components increase indefinitely as frequency is lowered. This means that the low frequency dispersion of the Bi_{48-x}Sb_xSe₅₂ alloys arise mainly from the localized charge carrier conduction, which is characterized by free electron transport mechanism. It is seen also that ε' and then decrease for high concentration of Sb content. This behavior is related to change of the free carrier concentration upon change the antimony concentration [7]. As mentioned above an increase of the free electron

concentration is dominated for low antimony content, whereas it is suppressed



Fig. (4): Real part of dielectric constant vs. frequency for Bi_{48-x}Sb_xSe₅₂ alloys.



Fig.(5): Imaginary part of dielectric constant vs frequency for Bi_{48-x}Sb_xSe₅₂ alloys

3.3. Optical properties:

3.3.1. Effect of Composition on the Absorption of IR Radiation:

Figure (6) illustrates the spectral behavior of the observed IR transmittance of the polycrystalline $Bi_{48-x}Sb_xSe_{52}$ alloys, where x = 0.6, 11, 25, 35 and 44 at. % in the wavelength range 2.5-25 μ m. It is shown that the general behavior of the spectra depends on the composition.



Fig. (6): IR Spectral distribution of $T(\lambda)$ for crystalline $Bi_{48-x}Sb_xSe_{52}$ alloys.

Figure (7) shows the variation of the absorption coefficient α as a function of photon energy (hv) for the Bi_{48-x}Sb_xSe₅₂ alloys. It is seen that, in the photon energy range 0.1 - 0.3 eV, α almost increases as the Sb content increases.



Fig. (7): Variation of the absorption coefficient (α) of Bi_{48-x}Sb_xSe₅₂ alloys with the photon energy(hv).

Below a certain photon energy (about 0.2 eV), α decreases sharply for all compositions of the system. It is known that near the fundamental absorption edge, the absorption coefficient varies rapidly with frequency.

The shape of the long-wavelength edge gives information on the scattering mechanism of the free charge carriers. For the propose to determine the values of the exponent **b** in the relation: $\alpha \sim \lambda^b$, the plot of $ln \alpha$ versus $ln \lambda$ shown in Fig. (8) gives b = 0.63 for x = 0.6 at.% which is less than 1.5 and the exponent b vanishes for other compositions with higher values of Sb content. It has been reported that in $(Bi_{1-x}Sb_x)_2Se_3$ crystals (x = 0 to 0.2) at temperature close to 300 K the dominant scattering mechanism of the free carriers is the scattering by acoustic branches of lattice vibrations [6]. This type of scattering takes place for only one composition (x = 0.6 at.%) of the Bi_{48-x}Sb_xSe₅₂ system.



Fig. (8): Plots of $\ln \alpha$ vs. $\ln \lambda$ for polycrystalline Bi_{48-x}Sb_xSe₅₂ alloys.

3.3.2. Compositional dependence of the optical energy gap

The spectral dependence of the transmittance (*T*) of the polycrystalline $Bi_{48-x}Sb_xSe_{52}$ alloys was measured in the wavelength range 0.32-1.1 μ m with normal incidence of light at room temperature. The absorption in semiconductors is characterized by Tauc`s relation:

$$\alpha h \upsilon = A (h \upsilon - E_g)^n \tag{5}$$

where A is a constant, α is the absorption coefficient, E_g is the optical energy gap of the studied materials and the exponent n characterize the transition type. This is the usual method to determine exactly the transition type. Fig. (9) shows the plots of $(\alpha h\nu)^2$ versus h ν for the Bi_{48-x}Sb_xSe₅₂ alloys (where x = 0.6, 11, 25, 35 and 44 at.%) indicating that the transition is allowed direct transition.



Fig. (9): Plots of $(\alpha h\nu)^2$ versus hv for the system Bi_{48-x}Sb_xSe₅₂.

The dependence of the optical energy gap Eg of the Bi_{48-x}Sb_xSe₅₂ alloys on composition x is shown in Fig. (10). There are two factors affecting the optical energy gap. One of them is related to the Moss – Burstein effect in which an increase in the free carrier concentration N leads to a shift of the edge towards higher energies whereas a decrease in the value of N give rise to a shift of the absorption edge towards lower energies. The second factor is the variation of the ionicity of the crystal which showed that the substitution of Sb atoms for Bi in the Bi₂Se₃ crystal lattice results in an observable increase of the optical gap E_g [6] (where the value of E_g for Bi_2Se_3 is 0.115 eV [8] and for Sb_2Se_3 is 1.2 eV [9]). Our results reveal that the action of ionicity effect is dominant. In samples with lowest Sb content (x = 0.6 and 11 at.%) the free carrier concentration increases as discussed above (section 3.2.1) but the energy gap was found to decrease where as in the samples with high Sb content (x = 25, 35 and 44 at. %) the optical gap is increased with rising the Sb content due to the increase in the ionicity. In general the variation of the activation energy of conduction with the composition is similar with that of the optical energy gap. This variation is related to change of the free carrier concentration as well as the variation of the ionicity.



Fig. (10): The variation of the optical energy gap E_g as a function of composition Bi_{48-x}Sb_xSe₅₂ alloys

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