

A constant force function for the lattice Boltzmann equation in isotropic turbulence

Waleed Abdel Kareem^{1,2}

¹ Suez University, Faculty of Science, Department of mathematics,

² Academy of Scientific Research & Technology, 101 Al Kasr El Aini, Cairo, Egypt.

ARTICLE INFO

Article history:

Received 21 February 2021

Received in revised form 17 April 2021

Accepted 17 April 2021

Available online 17 April 2021

Keywords

Lattice Boltzmann equation,
forcing function,
FORTRAN code.

ABSTRACT

A constant force function for the lattice Boltzmann method is introduced and parts of the programming code are presented. Adding force to the lattice Boltzmann equation is a challenging issue because the nature of the forcing techniques and the nature of the lattice Boltzmann equation (LBE) itself. The LBE is a difference equation and it is coded in the physical space, however, the forcing is usually done in the Fourier space at low wave-numbers. This paper discusses the force function and its coding to be readable by the LBE without losing of any physical meaning of the engineering problem.

1. Introduction

Simulation of forced isotropic turbulent flow is an important engineering problem because it is a standard problem to study statistical properties of turbulence and it is the best kind to investigate turbulent theories, visualizing vortices and test the Kolmogorov theories of turbulence. The forcing of the lattice Boltzmann equation is done by adding the force to the collision term or by shifting the velocity field or by adding the force to the collision process and shifting the velocity field at the same time. This process is done randomly each time step and the force is also added to the space at random points in the box each time step. In this study, the force function will be coded and all parameters will be discussed. The coding will be considered using the FORTRAN language and the code can be modeled easily with any other programming languages. Many studies have been presented to study the forcing of the lattice Boltzmann method such as Luo [1], Shan and Chen [2], Guo et al. [3], Succi [4], Cosgrove et al. [5], Siggia and Patterson [6], Mohseni et al. [7]. Generating the velocity field of isotropic turbulence using the forcing function was introduced by Abdel Kareem et al. [8] where extracting of multiscale vortical structures was investigated with resolution of 128^3 and the function is also used to resolution of 256^3 by comparing the lattice D3Q15 and D3Q19 models [9]. It was shown that the D3Q19 model is more stable than the D3Q15 model. Some recent studies of isotropic turbulent

flows are presented using the same forcing function by Albernaz et al.[10,11].

Elghobashi[12] applied the same force at low wave-numbers at every time step to generate a statistically stationary velocity field for Reynolds number between 73 and 133. Gkoudesnes and Deiterding[13] use the same forcing function to their LBM solver for homogeneous isotropic turbulence(HIT) in a periodic box and they started with zero initial velocity and unit density. A comparison between the direct numerical simulations (DNS) and large eddy simulations (LES)[14] is also considered using the same forcing function where they use the forcing scheme in their HIT for different resolutions starting from 32^3 to 512^3 . Also, Abdel kareem et al.[15] applied the force function in their study of filtering isotropic turbulent data in addition to synthetic data that are generated by solutions of Navier-Stokes equations. This paper is organized as follows: section 1 is the current introduction and section 2 presents the forcing function and its coding in FORTRAN language. In section 3, the results and discussion which are mainly concentrated in depicting the energy spectrum and the vortical structures resulted from the forcing function and finally, sec.4 is the conclusion of the study.

2. The forcing function and FORTRAN code

The lattice Boltzmann equation (LBE) can be written as

$$f_{\alpha}(\mathbf{x} + \mathbf{e}_{\alpha}\delta t, t + \delta t) - f_{\alpha}(\mathbf{x}, t) = \frac{-1}{\tau}(f_{\alpha}(\mathbf{x}, t) - f_{\alpha}^{eq}(\mathbf{x}, t)) + 3\rho w_{\alpha}(\mathbf{e}_{\alpha} \cdot \mathbf{F}).$$

Where the last term represents the force term which is added to the collision term in the lattice Boltzmann method (LBM). The basic motivation of this study is to present a

* Corresponding authors at: Suez University

E-mail addresses: waleed.abdelkareem@sci.suezuni.edu.eg
(W. Abdel Kareem)

function F that appeared in the last term in the equation. All the parameters and definitions of the functions introduced in the above LBE can be found in references [8,9], respectively. Also, the force code will be discussed and some parts will be introduced. The most important point is the function itself which should be clear in containing the coordinate variables and the wavenumbers. The function used can be written as:

$$F_1 = \begin{cases} F_x = A \sum_{k_1=-k_{max}}^{k_{max}} \sum_{k_2=-k_{max}}^{k_{max}} \sum_{k_3=-k_{max}}^{k_{max}} [(\frac{k_1^2}{k^2}) \sin(\frac{2\pi}{N} (k_1x + k_2y + k_3z) + \phi)] \\ F_y = -2A \sum_{k_1=-k_{max}}^{k_{max}} \sum_{k_2=-k_{max}}^{k_{max}} \sum_{k_3=-k_{max}}^{k_{max}} [(\frac{k_2^2}{k^2}) \sin(\frac{2\pi}{N} (k_1x + k_2y + k_3z) + \phi)], \\ F_z = A \sum_{k_1=-k_{max}}^{k_{max}} \sum_{k_2=-k_{max}}^{k_{max}} \sum_{k_3=-k_{max}}^{k_{max}} [(\frac{k_3^2}{k^2}) \sin(\frac{2\pi}{N} (k_1x + k_2y + k_3z) + \phi)] \end{cases}$$

where, the random phase is represented by ϕ and the wavenumbers are k_1, k_2 and k_3 in the 3D cartesian directions and the disturbance is caused by the sinusoidal function. A is the force amplitude and can be chosen equal to 10^{-4} . This force function has been used in several studies in the last ten years such as [10,11,12,13,14]. The first part of the code is generating the random phase function ϕ by a system random number generating function $rand()$:

```
do k1=-k_max, k_max
do k2=-k_max, k_max
do k3=-k_max, k_max
  phi(k1,k2,k3)=2*pi*rand()
enddo
enddo
enddo
```

k_{max} is the maximum cutoff wavenumber. The second part of the code is the main body of the code where the sine function is calculated at the low wavenumbers - $k_{max} \leq |k| \leq k_{max}$:

```
do i=1,n_x
  dx=2*pi*REAL(i)/REAL(nx)
do j=1,n_y
  dy=2*pi*REAL(j)/REAL(ny)
do k=1,n_z
  dz=2*pi *REAL(k)/REAL(nz)
  u_x(i,j,k)=0.0
  u_y(i,j,k)=0.0
  u_z(i,j,k)=0.0
do k1=-k_max, k_max
do k2=-k_max, k_max
do k3=-k_max, k_max
index=REAL(k1*k1+k2*k2+k3*k3)
if(sqrt(index).gt.0.and.sqrt(index).le.k_max)then
F=sin(k1*dx+k2*dy+k3*dz+phi(k1,k2,k3))
u_x(i,j,k)=u_x(i,j,k)+1.0/index*k2*k3*F
u_y(i,j,k)=u_y(i,j,k)-2.0/index*k1*k3*F
u_z(i,j,k)=u_z(i,j,k)+1.0/index*k1*k2*F
endif
```

```
enddo
enddo
enddo
u_x(i,j,k)=A*u_x(i,j,k)
u_y(i,j,k)=A*u_y(i,j,k)
u_z(i,j,k)=A*u_z(i,j,k)
Write(file-number) u_x(i,j,k),u_y(i,j,k),u_z(i,j,k)
enddo
enddo
enddo
```

Here n_x, n_y and n_z are the box lengths in x, y and z -directions. The output velocities are the low-wavenumber velocity modes that are calculated in the sphere with a radius of k_{max} . In all cases, one can set k_{max} equals to 2 or 3 or 4 and the forcing amplitude as $A=10^{-4}$. The force is divergence free, where it is clear that $div F=0$.

3. Results and discussion

Figure 1 shows the energy spectrum for this force, where the cutoff wavenumber effect can be noticed in the figure. Figure 2 shows the vortical structures that are depicted using the method introduced by Abdel Kareem [16] whose definition is:

$$Q_s^S = [(Q_W^3 + Q_s^3) + (\Sigma^2 - R_s^2)]^{\frac{1}{3}},$$

Where Q_W represents the rotation tensor strength and Q_s represents the deformation tensor strength, Σ and R_s are the enstrophy production term and the strain rate production, respectively. The mathematical definitions of the tensors are:

$$Q_W = \frac{1}{2} \Omega_{ij} \Omega_{ij}, Q_s = -\frac{1}{2} S_{ij} S_{ij}, \Sigma = \omega_i S_{ij} \omega_j, R_s = -\frac{1}{3} S_{ij} S_{jk} S_{ki}$$

Where Ω_{ij} is the rotation tensor and S_{ij} is the strain tensor. After few time steps, this force can develop the simulations of the lattice Boltzmann equation and the spectrum can be transformed from low wavenumbers to higher wavenumbers as depicted in Figure 3. Also, the large vortices that are visualized in Figure 1 will be developed to thinner and longer vortices with time advancement and it is shown in Figure 4. The developed figures 3 & 4 indicate that the forcing function is a good candidate for isotropic turbulence and can generate steady and stationary isotropic turbulent flow data. This forcing is succeeded to generate 3D isotropic turbulent data with statistical results that are close to the results obtained by Fourier spectral numerical simulations of Navier-Stokes equations. In recent studies, the same function is compared with other different forces to generate 3D isotropic turbulent data [17] and also used in comparison of different forcing techniques of the LBM [18].

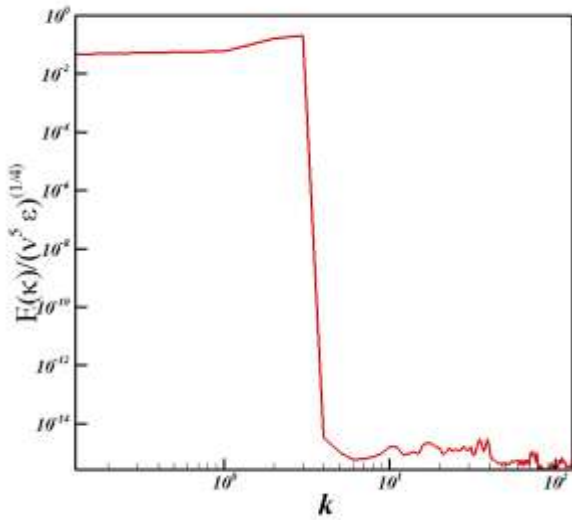


Figure 1: Energy spectrum of the forcing function at $n_x=n_y=n_z=256$.

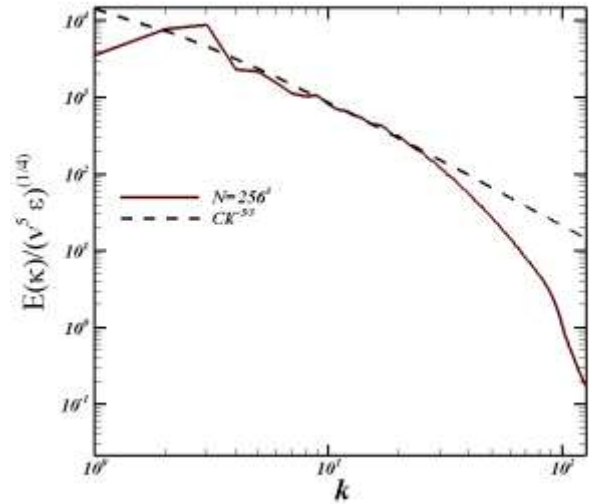


Figure 3: Energy spectrum of a developed time step generated using the forcing function at $n_x=n_y=n_z=256$ [9].

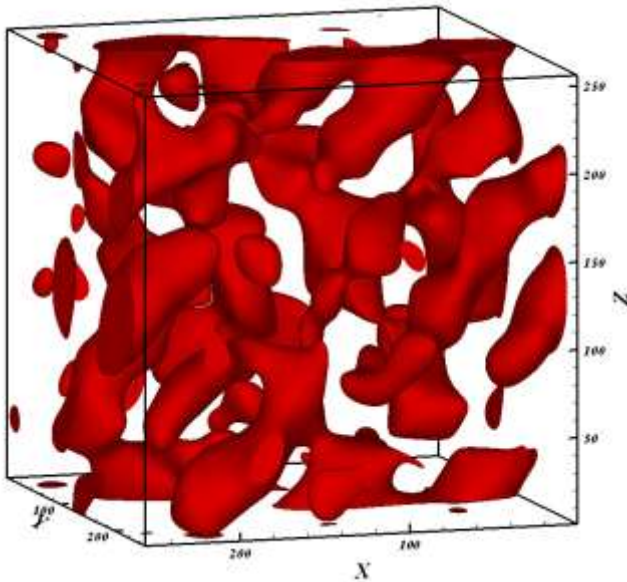


Figure 2: Isosurfaces of the vortices generated from the forcing function initial velocities at the first calculations with $n_x=n_y=n_z=256$.

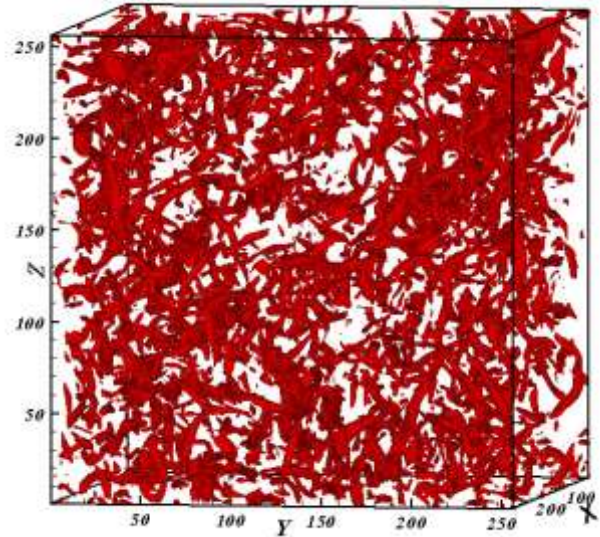


Figure 4: Isosurfaces of the vortices for a developed time step that is generated from the forcing function velocities at later time step with $n_x=n_y=n_z=256$ [9].

4. Conclusion

A force function is introduced to discuss the linear forcing of the lattice Boltzmann equation and the method can be also used with other numerical methods for the same purpose. The forcing function is divergence free and can compensate the turbulent velocity with adding energy and preventing the decaying of the flow field. The coding of the function using FORTRAN is also presented in a clear and a simple way to help the readers to adjust the programming of the function using any programming language. The constant force can be saved and then used at every time step of the simulations.

ASRT Acknowledgement

This paper was supported financially by the academy of Scientific Research and Technology (ASRT), Egypt, Grant Project no. 6477.

References

- [1] L. Luo, " Lattice gas automata and lattice Boltzmann equations for two dimensional hydrodynamics", PhD. thesis, Georgia Institute of Technology, (1993).
- [2] X. Shan and S. Chen, " Simulation of nonideal gases and gas-liquid phase transitions by the lattice Boltzmann equation, Phys. Rev. E 49(4),(1994):pp.2941-2948.

- [3] Guo Z.,Zheng C., Shi B., "Discrete lattice effects on the forcing term in the lattice Boltzmann method", Phys. Rev. E 65,2002(046308).
- [4] Succi S. The Lattice Boltzmann Equation for Fluid Dynamics and Beyond. Oxford University Press; 2001.
- [5] Cosgrove J., Buick J., Tonge S., Munro C., Greated C. and Campbell D.Application of the lattice Boltzmann method to transition in oscillatory flow. J. Physics A: Math. Gen. 2003; 36: pp.2609-2620.
- [6] Siggia E. and Patterson G. Intermittency effects in a numerical simulation of stationarythree-dimensional turbulence. J. Fluid Mech. 1987;86: pp.567-92.
- [7] Mohseni K., kosovic B., Shkoller S. and Marsden J. Numerical simulations of the Lagrangian averaged Navier-Stokes equations for homogeneous isotropic turbulence. Phys. Fluids 2003; 15:pp.524-44.
- [8] Abdel Kareem W., S. Izawa, A. Xiong and Y. Fukunishi, "Identification of Multi-scale Coherent Eddy Structures in a Homogeneous Isotropic Turbulence". Journal of Progress of Computational Fluid Dynamics, Vol. 6, (2006).
- [9] W. Abdel Kareem, S. Izawa, A. Xiong and Y. Fukunishi, " Lattice Boltzmann Simulations of Homogeneous Isotropic Turbulence", Computers & Mathematics with Applications: 58(2009) 1055-1061.
- [10] Albernaz D., Do-Quang M., Hermanson J. and Amberg G., Thermodynamics of a real fluid near the critical point in numerical simulations of isotropic turbulence. Physics of Fluids 28, (2016) 125105; doi:10.1063/1.4972276.
- [11] Albernaz D., Do-Quang M., Hermanson J. and Amberg G., Droplet deformation and heat transfer in isotropic turbulence. J. Fluid Mech.(2017), 820, pp. 61-85; doi:10.1017/jfm.2017.194.
- [12] Elghobashi S. Direct numerical simulations of turbulent flows laden with droplets or bubbles Annu. Rev. Fluid. Mech. 2019; 51: pp.217-244.
- [13] Gkoudesnes C., Deiterding R. Evaluating the LBM for large eddy simulations with dynamics sub-grid scale models 11th Int. Symposium on Turbulence and Shear Flow phenomena (TSFP11), Southampton, UK, July 30 to August 2, 2019.
- [14] Gkoudesnes C., Deiterding R. Verification and validation of a lattice Boltzmann method coupled with complex sub-grid scale turbulence models VI Int. conference on Particle-based Methods-Fundamentals and Applications,2019, E.Onate, M. Bischoff, D. Owen, R. Wriggers & T. Zohdi(Eds).
- [15] Abdel Kareem W., Izawa S., Klein M. and Fukunishi Y., A hyperbolic partial differential equation model for filtering turbulent flows. Computers & Fluids 2019; 190: pp.156-167.
- [16] Abdel Kareem W., A vortex identification method based on strain and enstrophy production invariants, Int. J. Mod. Phys. C, Vol. 31, No. 01, 2050003 (2020).
- [17] Abdel Kareem W., Djenidi L, Asker Z., Simulations of isotropic turbulence using lattice Boltzmann method with different forcing functions, Submitted 2021.
- [18] Abdel Kareem W., Djenidi L., Variable linear forcing techniques of the lattice Boltzmann method for turbulent flow simulations, submitted 2021