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MILITARY TECHNICAL COLLEGE

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CAIRO - EGYPT

### FULL -WAVE ANALYSIS OF THE GAP

## COUPLED MICROSTRIP ANTENNAS

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## ABSTRACT

Full-wave analysis of a gap is carried out using Galerkin's method in the Fourier transform domain (FTD). In the formulation of the problem, we consider two types of resonators. The analysis is carried our for both electric and magnetic walls at the planes of symmetry. The equivalent length  $\Delta l_e$  and  $\Delta l_m$  for electric and magnetic walls respectively are determined from the respective resonance frequencies of resonators. From  $\Delta l_e$  and  $\Delta l_m$  the discontinuity capacitances of the Pi- equivalent circuits are determined.

The equivalent circuit of a microstrip gap is calculated for different gap spacing.

# 1. INTRODUCTION

Applications of microstrip antenna are limited mainly because of their narrow bandwidth [1],[2]. A method for increasing the bandwidth of microstrip patch antenna by incorporating two additional resonators which are gap coupled to the radiating edges of a rectangular patch was developed by Gupta [3]. In this method [3], a two dimensional approach using impedance Green's function and the segmentation method has been used to analyze the proposed antenna configuration shown in Fig. 1-a, the antenna configuration is thus considered as a multiport network. The coupling gap between the two opened - rectangular microstrip patchs were modeled as two-dimensional capacitive pi-network as shown in Fig. 1-b. The values of these capacitances were obtained from the formulas for even and odd capacitances of coupled microstrips of unequal width [4]. Quasi-static characterization of capacitances with microstrip discontinuites are valid with sufficient accuracy, only up to few gigahertz. For a complete characterization of discontinuties, it is required to determined the frequency dependence of various parameters. This information is Obtained from full-wave analysis which is presented in this paper.

# 2. FORMULATION OF THE PROBLEM

Full-wave analysis of a gap between two antennas can be carried out by using Galerkin's method in FTD . In this case, one considers two types

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of resonators. First, an analysis is carried out with conducting planes (electric walls) at  $z = \pm D$  in the configuration shown in Fig. 2. A second computation is carried out for magnetic walls at  $z = \pm D$ . The dimension L is chosen such that (2D - 2L) is equal to the gap spacing,

#### Analysis of Microstrip Resonator with Two Electric Walls: 2.1.

As the fields existing in the structure shown in Fig. 2 can be considered as superpositions of TE (to z) and TM(to z) Fields, they can be expressed in terms of two types of scalar potentials  $\psi(x,y,z)$  and  $\varphi(x,y,z)$ . Accordingly, the fields existing in the structure are given by:

$$E_{zi} (x,y,z) = K_{i}^{2} \varphi_{i} + \frac{\partial^{2} \varphi_{i}}{\partial z^{2}}$$
(1a)

$$H_{zi} (x, y, z) = K_{i}^{2} \psi_{i} + \frac{\partial^{2} \psi_{i}}{\partial z^{2}}$$
(1b)

$$H_{xi} \quad (x,y,z) = j \omega \mathcal{E}_{i} \left( \frac{\partial \mathcal{V}_{i}}{\partial y} \right) + \frac{\partial \mathcal{V}_{i}}{\partial x \partial z} \tag{1c}$$

$$E_{xi} = (x, y, z) = \frac{\partial \varphi_i}{\partial x \partial z} - J \omega u_i \left(\frac{\partial \varphi_i}{\partial y}\right)$$
(1d)

where i=1,2 designates the substrate or the air region respectively  $K_1 = (\mathbf{\mathcal{E}}_r \mathbf{u}_r)^{\mathbf{V}_2} \mathbf{k}_0 \qquad \mathbf{\mathcal{E}}_1 = \mathbf{\mathcal{E}}_0 \cdot \mathbf{\mathcal{E}}_r \qquad \mathbf{u}_1 = \mathbf{u}_0 \mathbf{u}_r \qquad (2)$ (2a)  $K_2 = k_0 = \omega \left( \boldsymbol{\varepsilon}_0 \quad \boldsymbol{u}_0 \right)^{\boldsymbol{y_2}}, \boldsymbol{\varepsilon}_2 = \boldsymbol{\varepsilon}_0$ (2b) <sup>u</sup>2<sup>= u</sup>0

 $\omega$  is the operating frequency and  $\mathcal{E}_{o}$  and  $\mathcal{U}_{o}$  are the free-space permittivity and permeability, respectively. The boundary value problem associated with the structure is solved in the Fourier transform or the spectral domain. The original form of this method was first developed by Itoh and Mittra [6],[7]. Formulation of the problem for determining the resonant frequency is basically the same as that found by Itoh [5]. The essential difference here is that, discrete variable transform in z-direction k, is used in the formulation of the problem which is determined from the solutionof cavity resonator for instance.

$$k_{z} = (m - \frac{4}{2}) \frac{\pi}{D} \qquad \text{for } E_{x} \text{ even } - H_{x} \text{ odd } (\text{in } z) \text{ modes} \qquad (3a)$$
$$= m \frac{\pi}{D} \qquad \text{for } E_{x} \text{ odd } - H_{x} \text{ even } (\text{in } z) \text{ modes} \qquad (3b)$$

where m = 1, 2, ....

Following the same steps, which was developed in [5], the wave length  $\lambda_e$  corresponding to the resonant frequency  $f_{re}$  can be determined and  $\lambda_e$  is given by

> $\lambda_e = \frac{C}{f_{e} \sqrt{\epsilon_{e}}}$ (4a)

where

C

E<sub>re</sub>

is the velocity of light in free space

is the effective dielectric constant [8], where

$$\xi_{re} = \frac{\xi_{r} + 1}{2} + \frac{\xi_{r} - 1}{2} F(W/h)$$
 (4b)

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+  $0.04(1-2 \text{ W/h})^2$ +12h/2W)-1/2  $2W/h \leq 1$ F(W/h) =(4c)1+12h/2w)-1/2 2W/h≥1

# 2.2 Analysis of Microstrip Resonator with Two Magnetic Walls:

Following the same procedures stated in section 2.1, after replacing the two electric walls by two magnetic walls, and the only difference will be in the value of  $K_{z}$ where,

$$K_{z} = m \pi/D \quad \text{for } E_{x} \text{even} - H_{x} \text{ odd (in z) modes}$$
(5a)  
=  $(m - \frac{1}{2}) \frac{\pi}{D} \text{ for } E_{x} \text{ odd} - H_{y} \text{ even (in z) modes}$ (5b)

By the same way, the wave length  $\lambda$  m corresponding to the resonant frequency  $f_{\rm rm}$  of this structure can be determined and  $\lambda$ m is given by

$$\lambda_m = \frac{C}{f_{rm} \int \mathcal{E}_{re}}$$
(6)

# 2.3 Calculation of the Equivalent Circuit of the Gap:

The end effect of the open-ended microstrip patch can be represented by an extension in the length  $\Delta l_e$  and  $\Delta l_m$  for the two cases 2.1 and 2.2 respectively. tively, where

$$2(L + \Delta l_e) = \lambda e/2$$
(7a)

 $2(L + \Delta 1) = \lambda_m/2$ From Al and Al, two equivalent sapacitances C. and C [8] are

$$C_{e} = \frac{\int \mathcal{E}_{re}}{C Z_{o}} \Delta^{1}_{e}$$
(7b)  
and  $\Delta^{1}_{m}$ , two equivalent sapacitances  $C_{e}$  and  $C_{m}$ [8] are calculated  
(8a)

$$C_{m} = \frac{I \mathcal{E}_{re}}{C Z_{o}} \Delta^{1}_{m}$$
(8b)

$$Z_{0} = \frac{7}{2\pi\sqrt{\epsilon_{r_{e}}}} \ln\left(\frac{8h}{2W} + 0.25\frac{2W}{h}\right) \qquad (2W/h \le 1) \qquad (8c)$$

$$= \frac{\gamma}{\sqrt{\epsilon_{re}}} \left(\frac{2W}{h} + 1.393 + 0.667 \ln\left(\frac{2W}{h} + 1.444\right)\right)^{-1}$$
(8d)  
(2W/h≥1)

7 = 120 TT ohm where

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where,  $Z_0$  is the characteristic impedance of the microstrip which is evaluated using work of Wheeler and Schneider [6]. These capacitances  $C_e$  and  $C_m$  are related to the discontinuites of the Pi- equivalent circuit of a gap by the following equations

 $Cg = (C_{e} - C_{m})/2$ 

(9a) (9b) ٦

Cp= C<sub>m</sub>

# 3. RESULTS AND DISCUSSIONS

Numerical Computations have been Carried out for the resonant frequencies and the equivalent circuit of the gap Coupling

Fig. 3shows the resonant frequency  $\rm f_{re}$  and  $\rm f_{rm}$  versus the gap spacing. This Figure shows that as the spacing increases, the resonant frequency  $\rm f_{re}$  decreases,  $\rm f_{rm}$  increases.

The imaginary part in the computed resonant frequency is due to the energy Lost by radiation from the fringing fields at the open ends.

The equivalent circuit for the gap, as referred to the center Line of the gap, will include series capacitance and some negative shunt capacitance to account for the fact that the gap reduces the shunt capacitance in the vicinity of the center Line, O Liner [9].

Fig.4 shows the equivalent Capacitances Cg and Cp which decrease as the spacing increases. Table 1- Sumarizes the different results.

# 4. CONCLUSION

The Capacitive network model which was used [3] to represent the coupling between the main resonator and the additional resonators is Considered as very approximate . Thus, the full-wave analysis which is presented in this paper for determining the gap Coupling to improve the accuracy of representing the Coupling.

It is planned in the future to obtain the equivalent Circuit of the gap Coupling for radiating edge gap Coupled micro strip antennas (REGCMA). AR-4 1189

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Fig.1.(a)

 Radiating edge gap coupled microstrip antenna with different parasitic element

(b) Lumped element network for modeling of coupling gaps

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	JS	TABLE 1 Jmmary of the results (ve	ersus gap spacing)	×
Spacing (CM)	0.15	0.165	0.18	0.19
f <sub>re</sub> [GHZ]	4.84411-j.2506 × 10 <sup>-6</sup>	4.84553-j.2849 × 10 <sup>-6</sup>	4.84696-j.3301 x 10-6	4.84791-j.3690 × 10-6
f <sub>rm</sub> [GHZ]	4.99190-j.9158 x 10 <sup>-8</sup>	4.98628-j.1577 × 10 <sup>-7</sup>	4.98078-j.2276 × 10 <sup>-7</sup>	4.97716-j.27651×10 <sup>-7</sup>
λ <sub>e</sub> (CM)	3.9892	3.9880	3.9869	3.9861
λ <sub>m</sub> (CM)	3.8711	3.8755	3.8798	3.8826
Δl <sub>e</sub> (CM)	-0.35268	-0.35297	-0.35327	-0.35346
∆ 1 <sub>m</sub> (CM)	-0.38221	038112	-0.38005	-0.37934
Cg (PF)	0.08842	0.0842811	0.080201	0.077504
C <sub>p</sub> (PF)	-2.2893	-2.2828	-2.27639	-2.27217
These results a	are for the following stru	cture parameters		
1 = 1.35 cm	W = 1.9	95 cm		

d = 0.159Cm

ч,

= 2.55

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